

Mesoscopic Simulation of Flowing Topological Composite Materials

Oliver Henrich

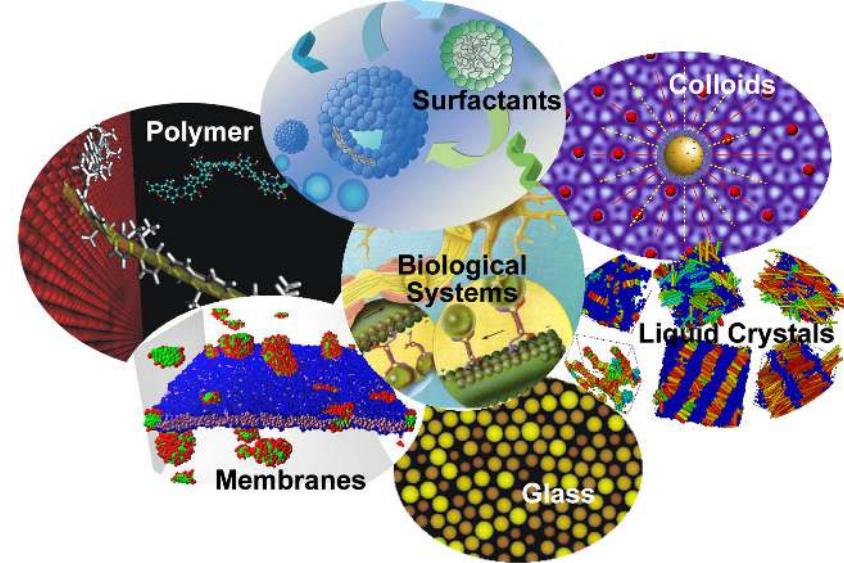
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Scottish Numerical Methods Network
Multiscale Computational Modelling

*International Centre for Mathematical Sciences
Edinburgh, 16th May 2018*



Outline

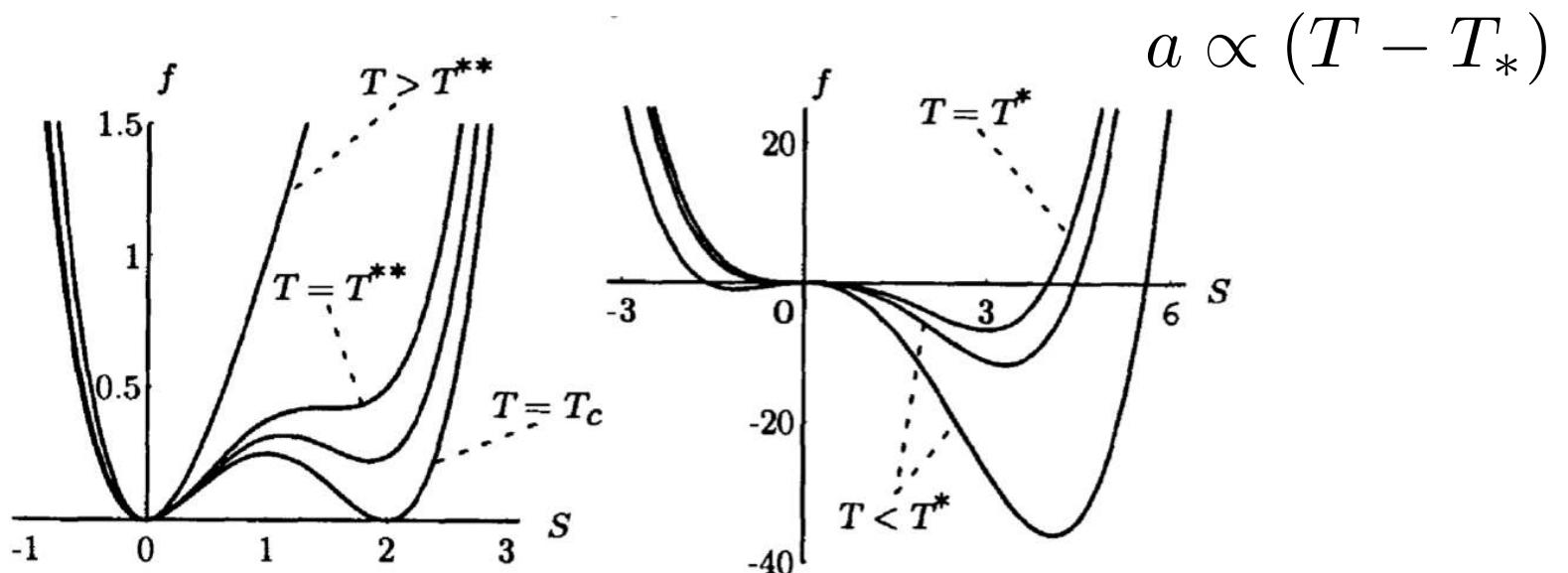


- Soft condensed matter
 - easily *deformed by thermal fluctuations*
 - *self-organisation* into *mesoscopic* structures
 - *macroscopic behaviour* depends on *mesoscopic structures*
- Landau models for phase transitions and complex fluids
- Lattice-Boltzmann simulation method
- Microfluidics of nanoparticle-liquid crystal composite materials



Landau Theory 1st Order Phase Transition

- **Scalar order parameter** $S \in \mathbb{R}$ measures "the order"
- **Free energy functional** $\mathcal{F}[S] = \int d^3x \{aS^2 - bS^3 + cS^4\}$



T^*, T^{**} limits of metastability for supercooling and superheating

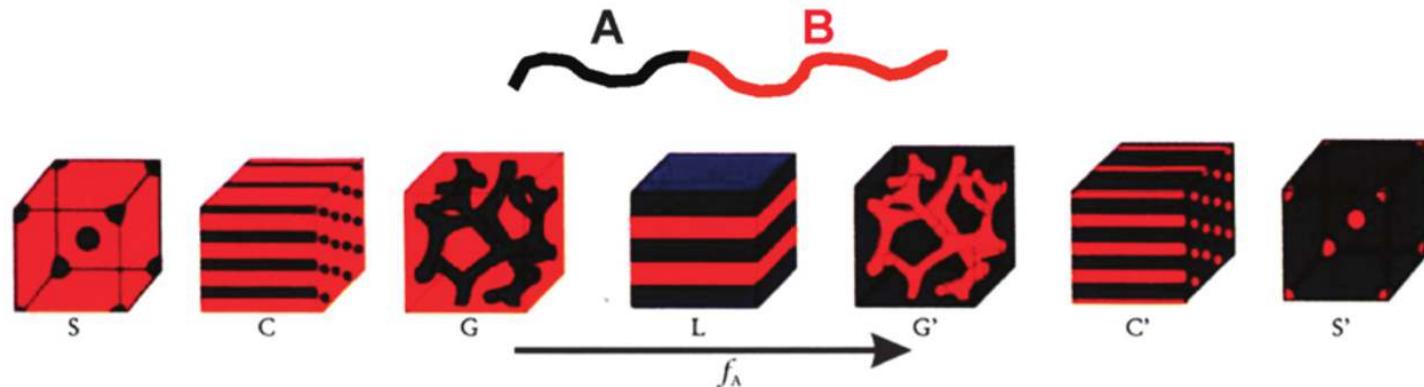


Discontinuous change of S at T_c



Binary Fluid: Order Parameter

- Phase-separating binary liquid: phase A, B, e.g. block-copolymers



- **Scalar order parameter** $\phi = \rho_A - \rho_B$ density difference between the two phases
- **Free energy functional**

$$\mathcal{F}[\phi] = \int d^3x \left\{ \frac{a}{2} \phi^2 + \frac{b}{4} \phi^4 + \frac{\kappa}{2} (\vec{\nabla} \phi)^2 \right\}$$

a, b Landau parameters
 κ related to surface tension



Binary Fluid: Dynamics

- **OP dynamics:** advection-diffusion equation

$$\partial_t \phi + \vec{\nabla} \cdot (\vec{u} \phi) = M \nabla^2 \frac{\delta \mathcal{F}}{\delta \phi} = M \nabla^2 \mu$$

μ chemical potential

M mobility coefficient

- **HD:** Navier-Stokes equation $\partial_t u_\alpha + (u_\beta \nabla_\beta) u_\alpha = \partial_\beta \sigma_{\alpha\beta}$

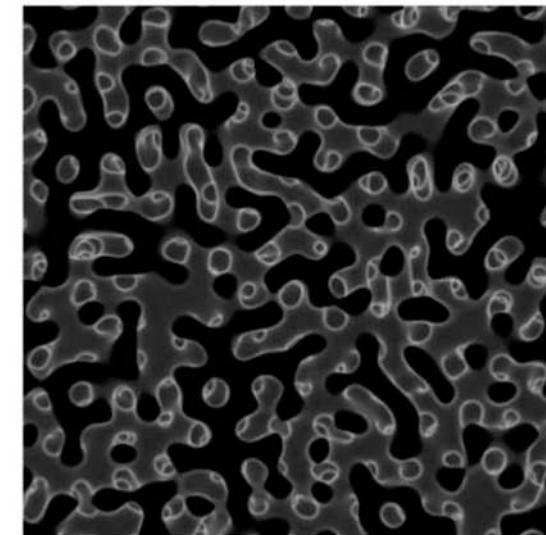
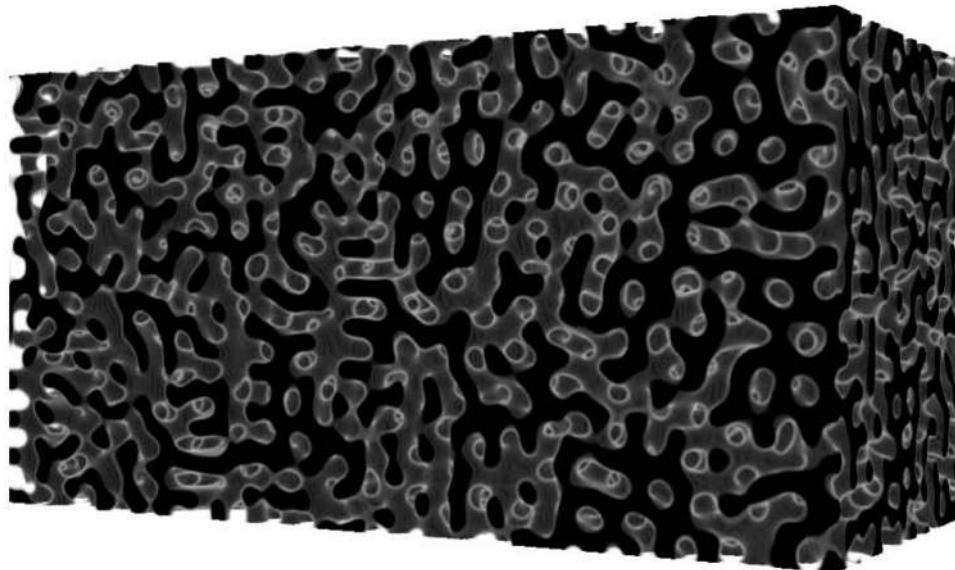
- **Stress tensor** in NSE

$$\sigma_{\alpha\beta} = -p \delta_{\alpha\beta} + \eta(\partial_\alpha u_\beta + \partial_\beta u_\alpha) - P_{\alpha\beta}^{chem}$$

$$P_{\alpha\beta}^{chem} = \kappa(\partial_\alpha \phi)(\partial_\beta \phi) \quad \text{chemical pressure due to composition}$$



Binary Fluid in Shear Flow

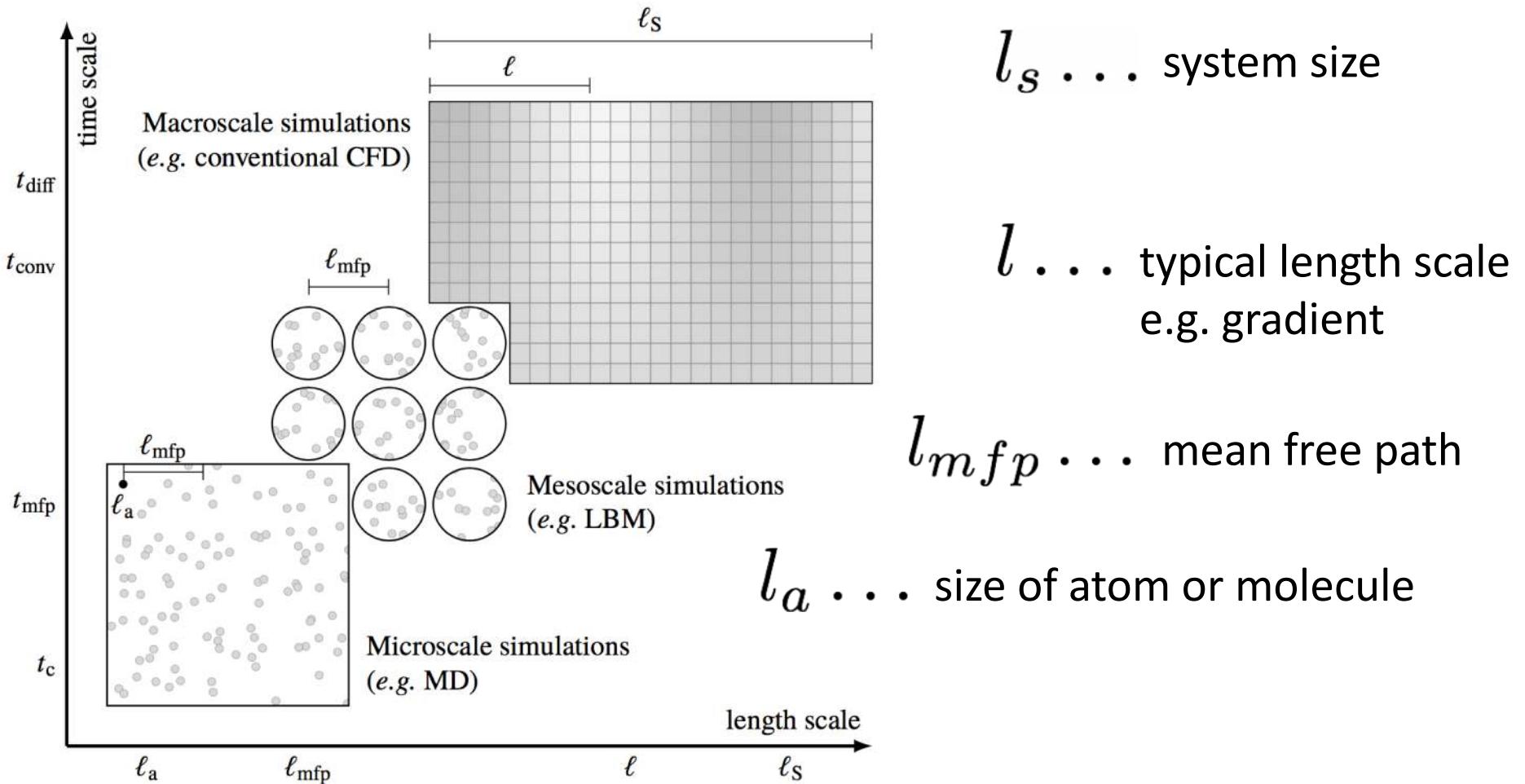


3D: undergoes ordering

2D: forms fluctuating
lamellae with spatial
heterogeneities



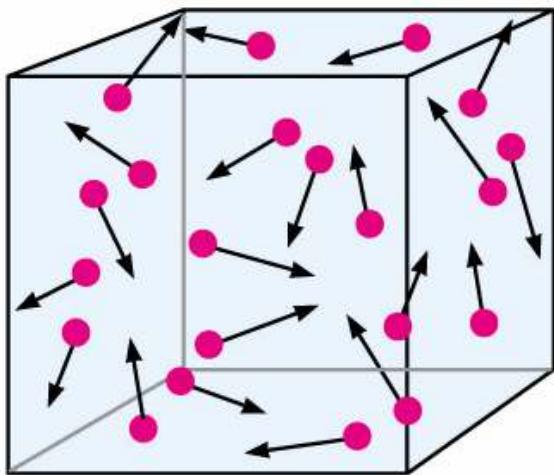
Mesoscopic Simulation Methods



Recent review: Ulf Schiller, Timm Krüger, Oliver Henrich,
Mesoscopic Modelling and Simulation of Soft Matter,
Soft Matter **14**, 9 (2018)



Lattice Boltzmann Method



Probability density function
 $f(\mathbf{r}, \mathbf{p}, t)$

Kinetic Theory: Boltzmann Equation
 $\partial_t f + \mathbf{u} \cdot \partial_{\mathbf{r}} f + \mathbf{F} \cdot \partial_{\mathbf{p}} f = \Omega[f]$

Density

$$\rho(\mathbf{r}, t) = \int f(\mathbf{r}, \boldsymbol{\xi}, t) d^3 \boldsymbol{\xi}$$

Momentum

$$\rho(\mathbf{r}, t) \mathbf{u}(\mathbf{r}, t) = \int \boldsymbol{\xi} f(\mathbf{r}, \boldsymbol{\xi}, t) d^3 \boldsymbol{\xi}$$

Energy

$$\rho(\mathbf{r}, t) e(\mathbf{r}, t) = \int |\boldsymbol{\xi} - \mathbf{u}(\mathbf{r}, t)|^2 f(\mathbf{r}, \boldsymbol{\xi}, t) d^3 \boldsymbol{\xi}$$

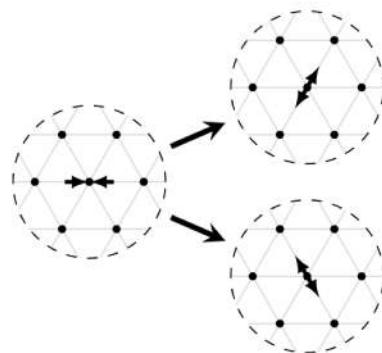


Lattice Boltzmann Method

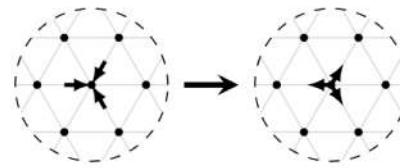
Discretisation of time, configuration AND momentum space

Lattice Gas Cellular Automata

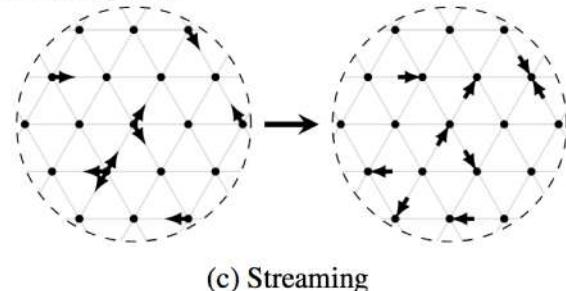
simplified collision rules reproduce HD



(a) Two-particle collision; the resolution is chosen randomly from the two options



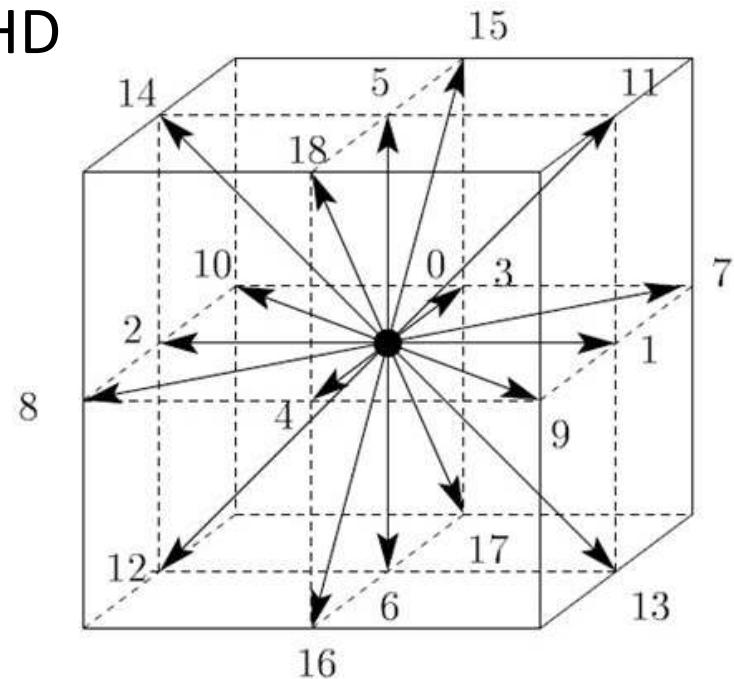
(b) Three-particle collision



(c) Streaming



U.Frisch, B.Hasslacher, Y.Pomeau,
Phys. Rev. Lett. **56**, 1505 (1986)



Discrete lattice velocities c_i

D3Q19: $i = \{0, \dots, 18\}$



Lattice Boltzmann Method

Lattice Boltzmann equation: for every discrete velocity \mathbf{c}_i

$$f_i(\mathbf{r} + \mathbf{c}_i, t + \Delta t) = f_i(\mathbf{r}, t) + \Omega_i[\{f, f_i^{eq}\}]$$

Continuous analogue of

$$\partial_t f + \mathbf{u} \cdot \partial_{\mathbf{r}} f + \mathbf{F} \cdot \partial_{\mathbf{p}} f = \Omega[f]$$

Collision operator

$$\Omega[\{f_i, f_i^{eq}\}] = \frac{1}{\tau_f} (f_i^{eq} - f_i)$$

$$\tau_f \propto \nu \quad \text{viscosity}$$

P. Bhatnagar, E. Gross , M. Krook,
“A model for collision processes in
gases”, *Phys. Rev.* **94**, 511 (1954)

Equilibrium distribution

$$f_i^{eq} = \omega_i \rho \left(1 + \frac{\mathbf{c}_i \cdot \mathbf{u}}{c_s^2} + \frac{(\mathbf{c}_i \cdot \mathbf{u})^2}{2c_s^4} - \frac{\mathbf{u}^2}{2c_s^2} \right)$$



Lattice Boltzmann Method

1. Collision step $f_i^*(\mathbf{r}, t) = f_i(\mathbf{r}, t) + \Omega_i[\{f, f_i^{eq}\}]$

2. Streaming step $f_i(\mathbf{r} + \mathbf{c}_i, t + \Delta t) = f_i^*(\mathbf{r}, t)$

or combined $f_i(\mathbf{r} + \mathbf{c}_i, t + \Delta t) = f_i(\mathbf{r}, t) + \Omega_i[\{f, f_i^{eq}\}]$

From **distributions** $f_i(\mathbf{r}, t)$ to **macroscopic observables**

Density $\rho = \sum_i f_i$

Momentum $\rho u_\alpha = \sum_i c_{i,\alpha} f_i$

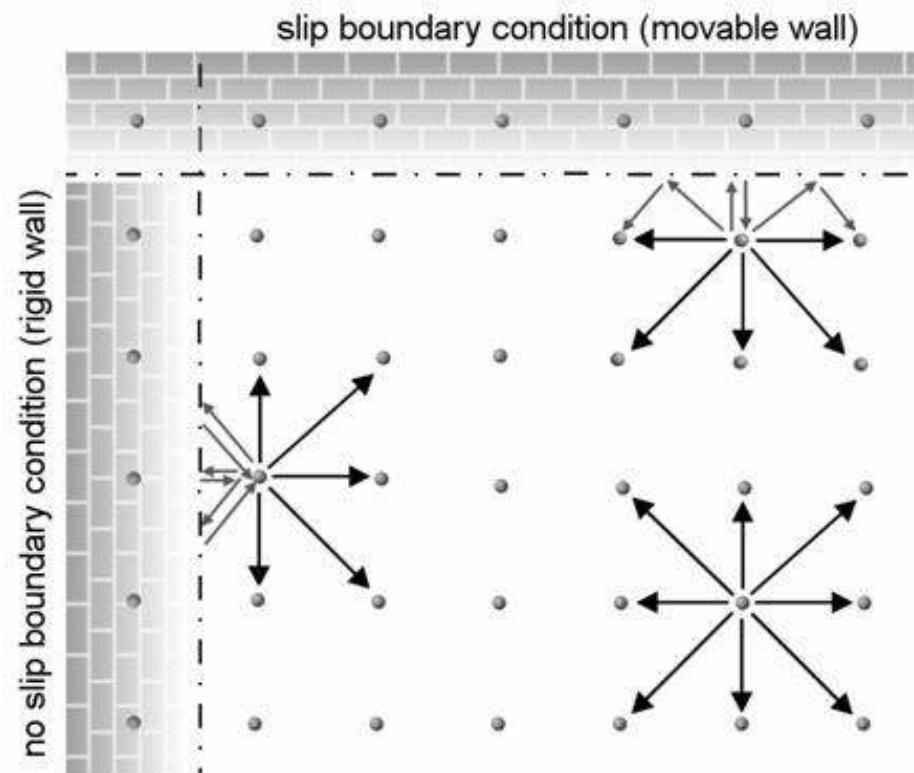
Stress $\sigma_{\alpha\beta} = \sum_i c_{i,\alpha} c_{i,\beta} f_i$



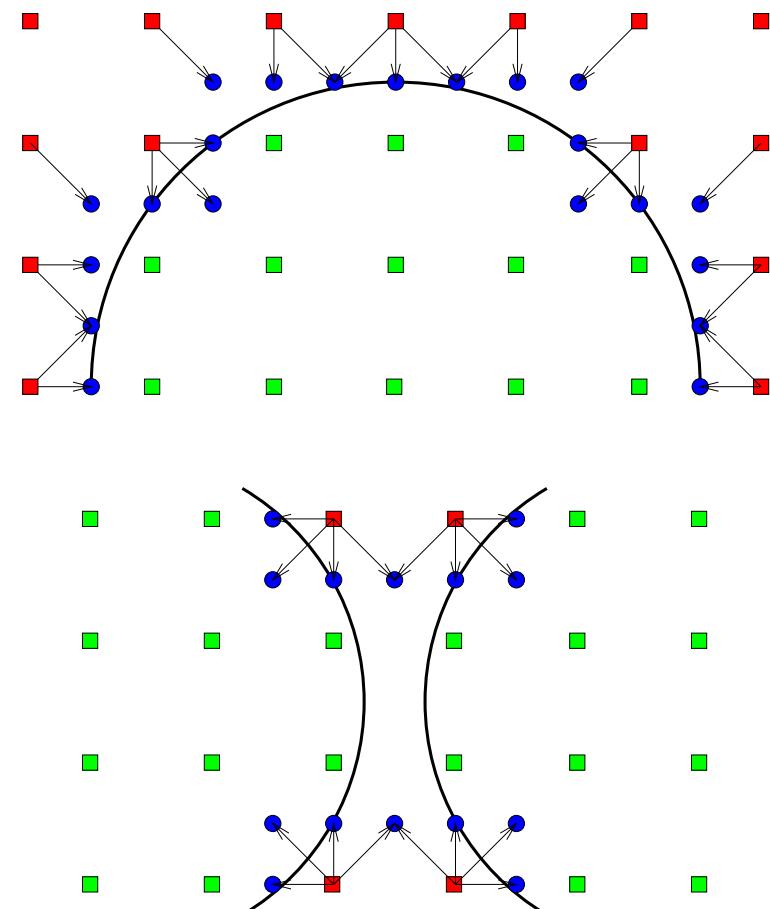
Lattice Boltzmann Method

(Relatively) simple boundary conditions

Walls

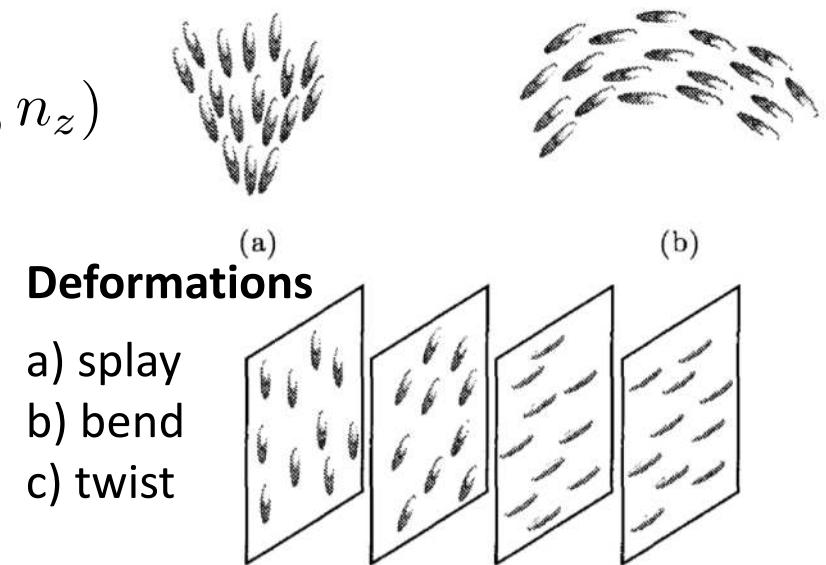


Particles



Mesoscopic Description of LCs

- **Unit vector**, aka ‘director’ $\vec{n} = (n_x, n_y, n_z)$
- **Scalar order parameter** $S \in [0, 1]$
- **Tensor order parameter** $Q_{\alpha\beta} = S \left(n_\alpha n_\beta - \frac{1}{3} \delta_{\alpha\beta} \right)$



Landau-de Gennes Free Energy Functional

$$\mathcal{F} = \int d^3r \left(\frac{A_0}{2} \left(1 - \frac{\gamma}{3} \right) Q_{\alpha\beta}^2 - \frac{A_0\gamma}{3} Q_{\alpha\beta} Q_{\beta\gamma} Q_{\gamma\alpha} + \frac{A_0\gamma}{4} (Q_{\alpha\beta}^2)^2 \right. \\ \left. + \frac{K}{2} \left\{ (\partial_\alpha Q_{\alpha\beta})^2 + (\varepsilon_{\alpha\mu\nu} \partial_\mu Q_{\nu\beta} + 2q_0 Q_{\alpha\beta})^2 \right\} - \frac{\epsilon_a}{12\pi} E_\alpha Q_{\alpha\beta} E_\beta \right)$$

Bulk terms

$A_0 \dots$ Bulk free energy constant

Elasticity terms

$\gamma \dots$ Temperature

Coupling to external field

$K \dots$ Elastic constant



Mesoscopic Description of LCs

EOM of the Q-tensor order parameter (finite difference)

$$(\partial_t + u_\alpha \partial_\alpha) \mathbf{Q} - \mathbf{S}(\{\partial_\beta u_\gamma\}, \mathbf{Q}) = \Gamma \mathbf{H}$$

Advection	Order-flow coupling	Molecular field
Molecular field	$\mathbf{H} = -\frac{\delta \mathcal{F}}{\delta \mathbf{Q}} + \mathbb{1}/3 Tr \frac{\delta \mathcal{F}}{\delta \mathbf{Q}}$	Rotational diffusion constant Γ
Order-flow coupling		
$\mathbf{S}(\mathbf{W}, \mathbf{Q}) = (\xi \mathbf{A} + \boldsymbol{\Omega})(\mathbf{Q} + \frac{\mathbb{1}}{3}) + (\mathbf{Q} + \frac{\mathbb{1}}{3})(\xi \mathbf{A} - \boldsymbol{\Omega}) - 2\xi(\mathbf{Q} + \frac{\mathbb{1}}{3}) Tr(\mathbf{Q} \mathbf{W})$		

Navier-Stokes & continuity equation (lattice-Boltzmann)

$$\rho (\partial_t u_\alpha + u_\beta \partial_\beta u_\alpha) = \partial_\beta (\sigma_{\alpha\beta}^s + \sigma_{\alpha\beta}^a) \quad \partial_\alpha u_\alpha = 0$$

$$\begin{aligned} \sigma_{\alpha\beta}^s &= -P_0 \delta_{\alpha\beta} + \eta \partial_\beta \{\partial_\alpha u_\beta + \partial_\beta u_\alpha\} - \xi H_{\alpha\gamma} \left(Q_{\gamma\beta} + \frac{1}{3} \delta_{\gamma\beta} \right) \quad \sigma_{\alpha\beta}^a = Q_{\alpha\gamma} H_{\gamma\beta} - H_{\alpha\gamma} Q_{\gamma\beta} \\ &\quad - \xi \left(Q_{\alpha\gamma} + \frac{1}{3} \delta_{\alpha\gamma} \right) H_{\gamma\beta} + 2\xi \left(Q_{\alpha\beta} + \frac{1}{3} \delta_{\alpha\beta} \right) Q_{\gamma\nu} H_{\gamma\nu} \\ &\quad - \partial_\alpha Q_{\gamma\nu} \frac{\delta \mathcal{F}}{\delta \partial_\beta Q_{\gamma\nu}} \end{aligned}$$

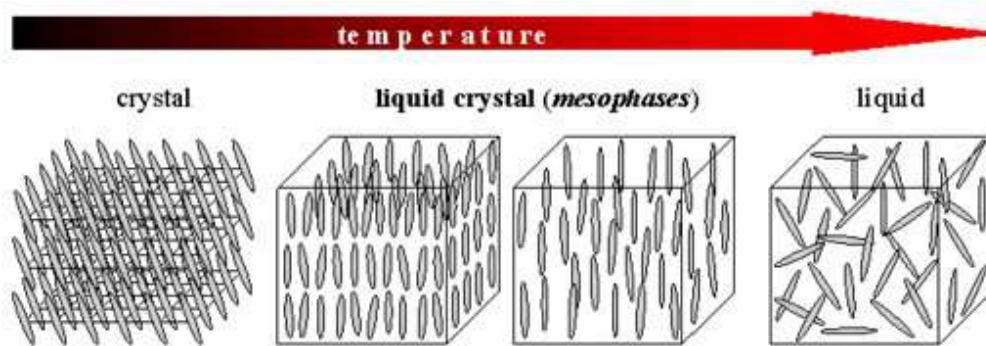


Liquid Crystal Mesophases

Solid, liquid crystal, liquid

Nematic

Smectic A & C



- 3-D lattice
- orientation
- solid

↳ *anisotropic*

- 1- (2-)D lattice
- orientation
- fluid

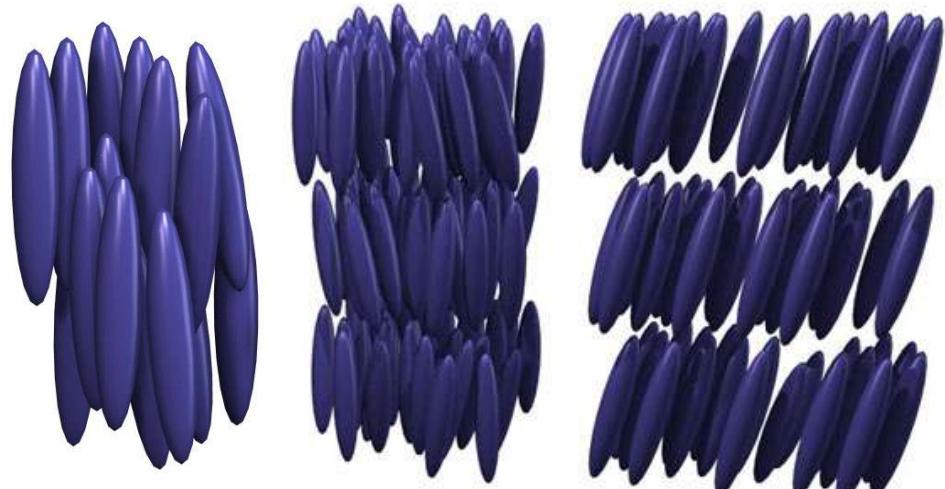
↳ *anisotropic*

- no lattice
- orientation
- fluid

↳ *anisotropic*

- no lattice
- no orientation
- fluid

↳ *isotropic*



Smectic A

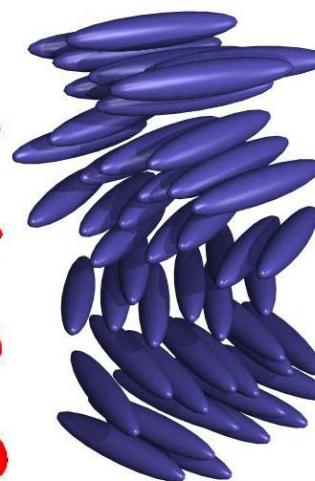
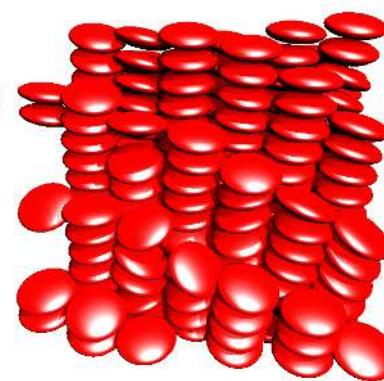
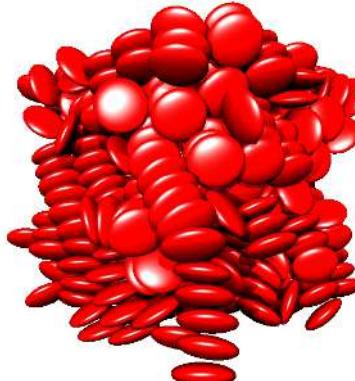
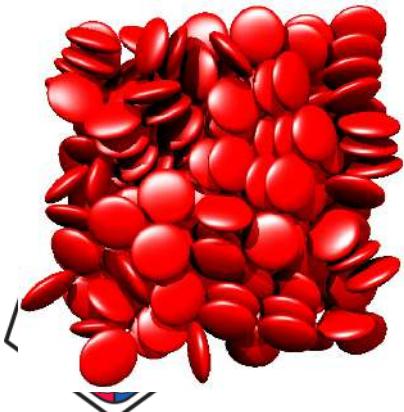
Smectic C

Isotropic

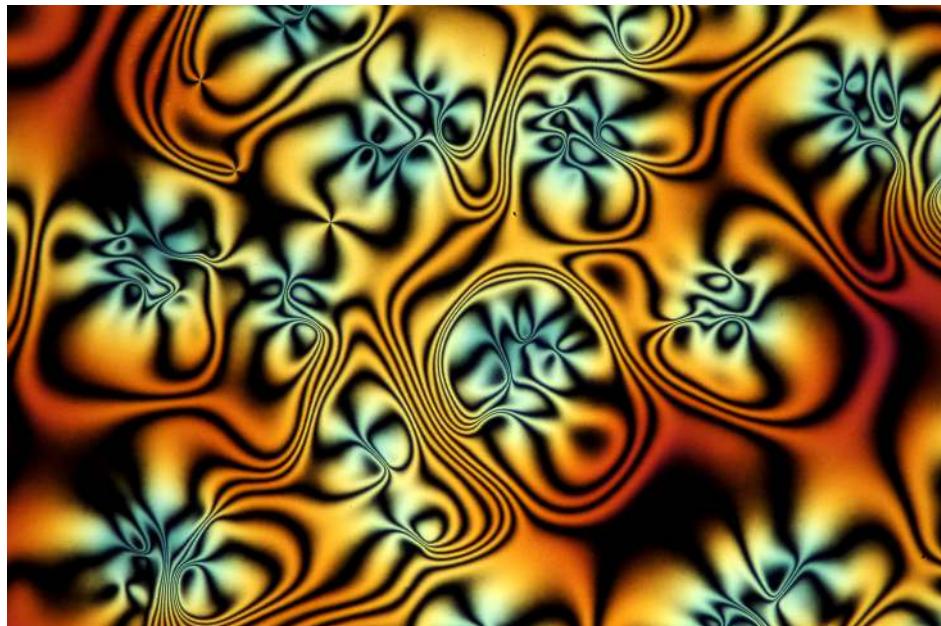
Nematic

Columnar

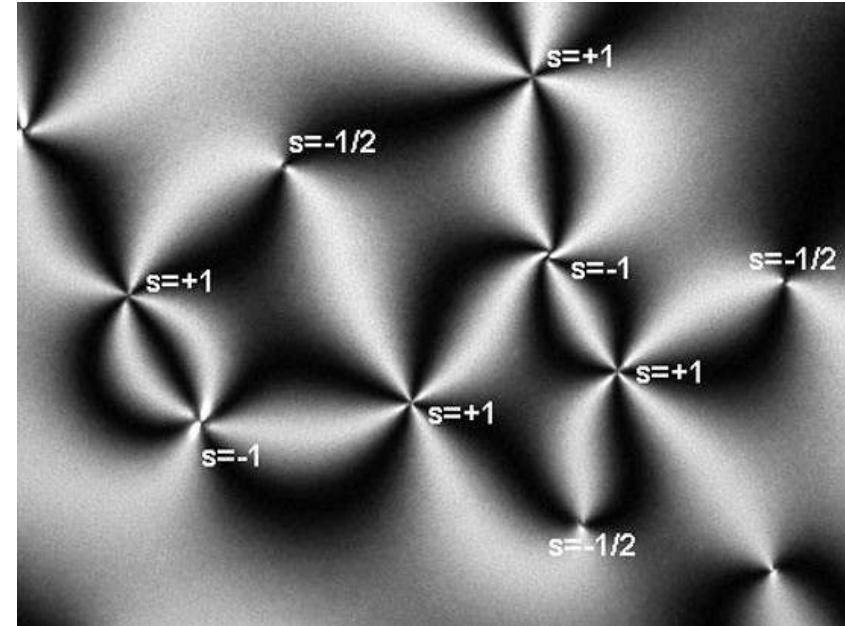
Cholesteric



Liquid Crystals: Topological Defects



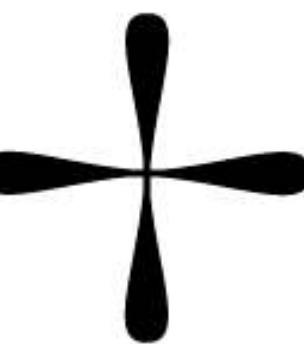
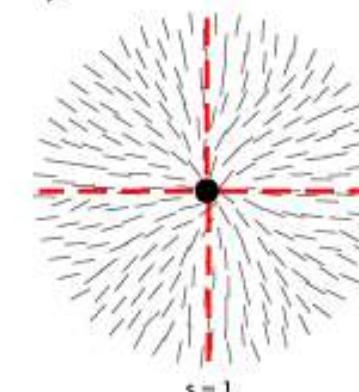
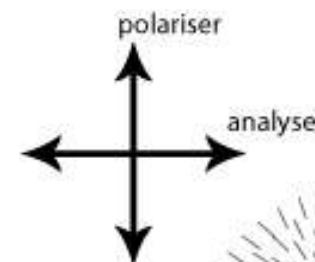
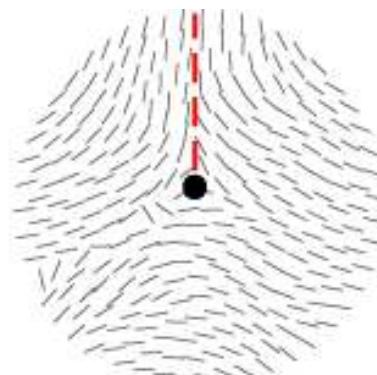
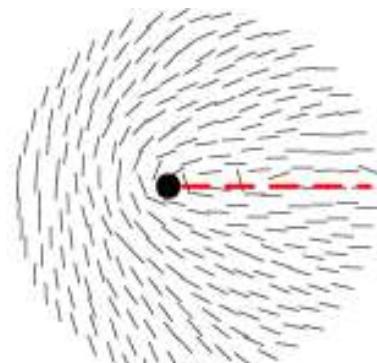
'Schlieren' texture in nematic
liquid crystals under **crossed**
polarisers



Defects carry a
topological charge



Liquid Crystals: Topological Defects



**Local order structure
and appearance in
polarised light
microscopy**

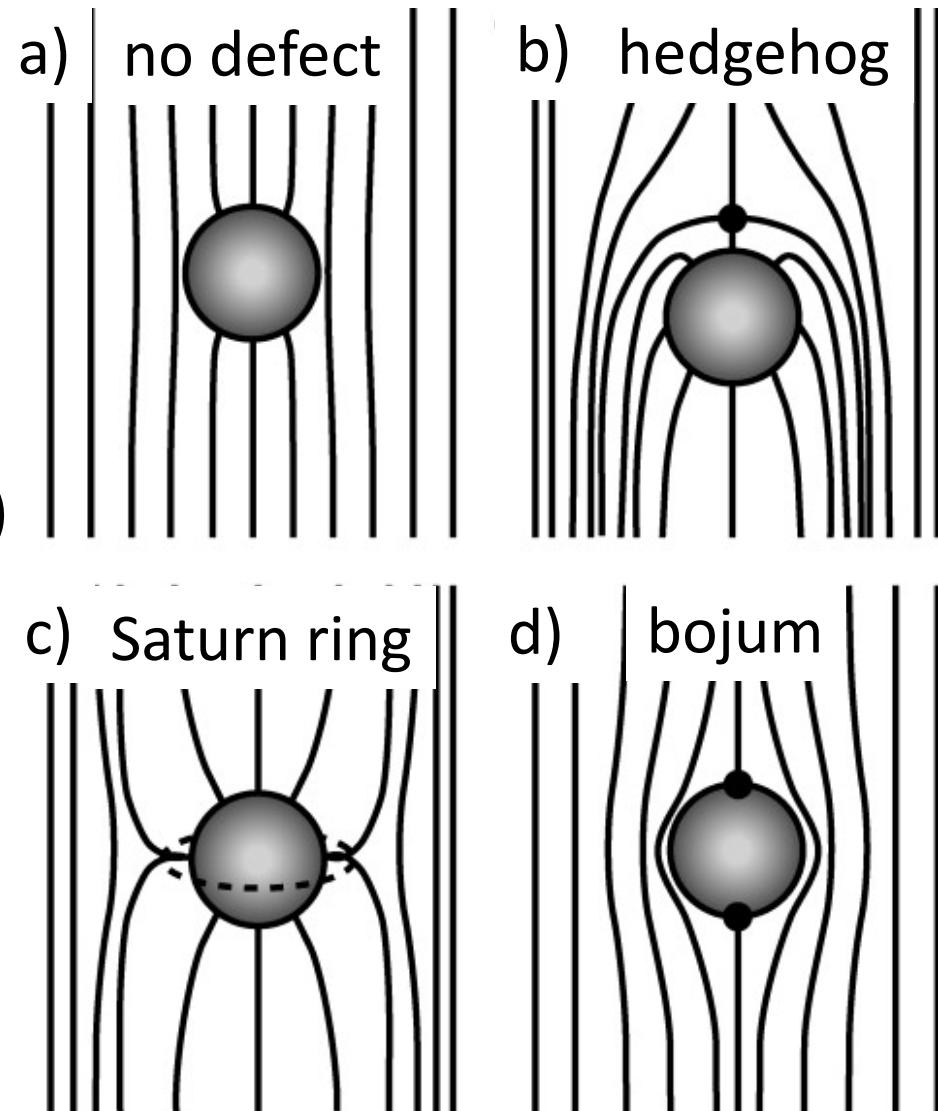


Liquid Crystals and Colloids

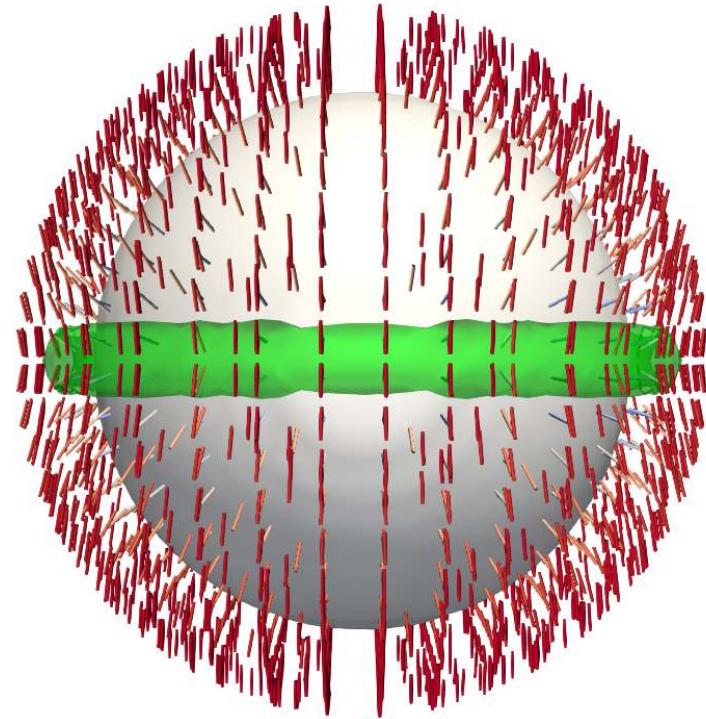
- Anchoring (preferred orientation) of the LC on the particle surface
- Free energy functional

$$\mathcal{F} = \int d^3r (f_{LdG} + f_s)$$
$$f_s = w(Q_{\alpha\beta} - Q_{\alpha\beta}^0)^2$$

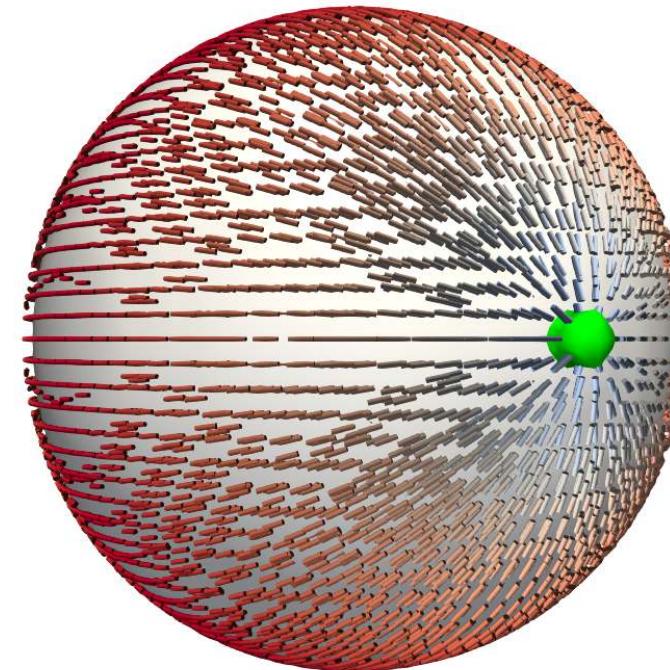
- a) weak normal
- b) very strong normal
- c) normal
- d) planar degenerate



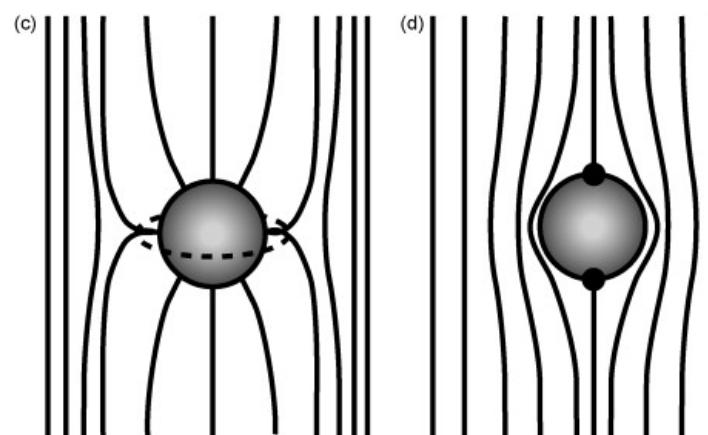
Liquid Crystals and Colloids



Normal anchoring
Saturn ring defect

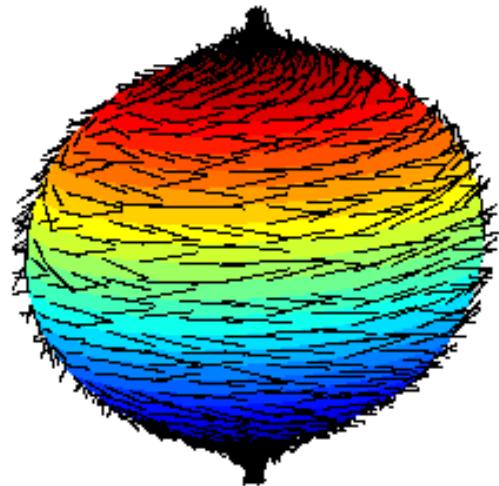


**Planar degenerate
anchoring**
bojum defect

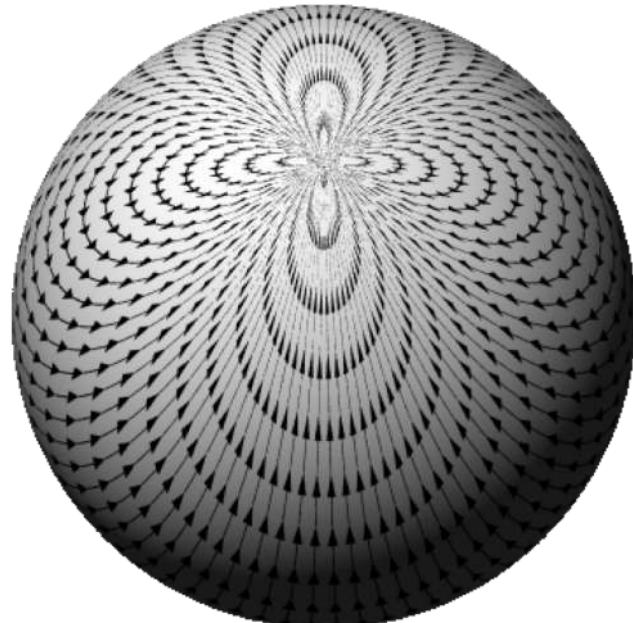
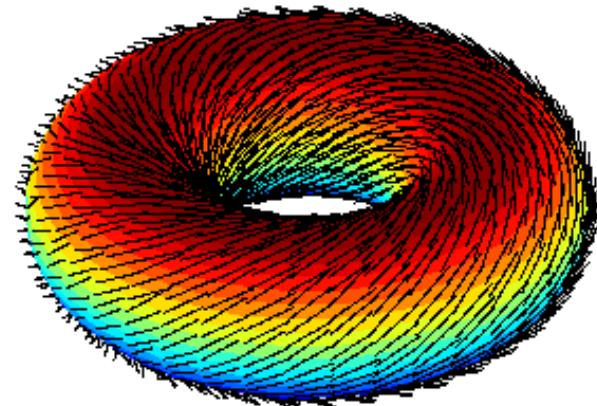


Liquid Crystals and Colloids

Hairy ball or hedgehog theorem (*Henri Poincaré, late 19th century*)

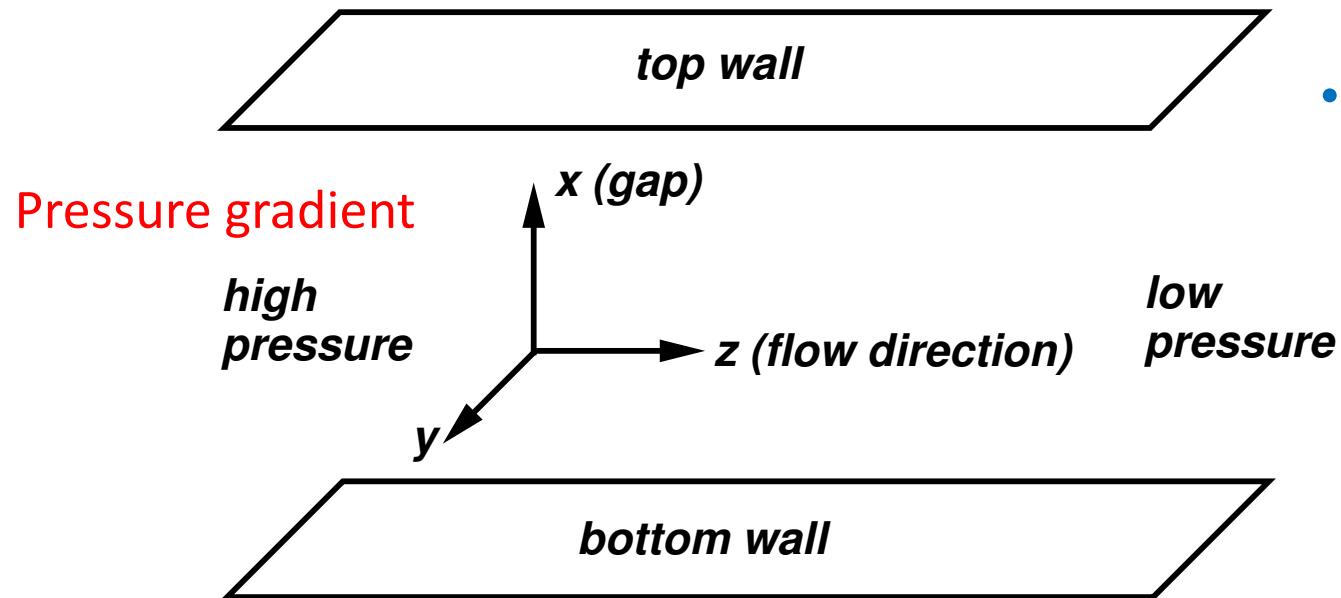


no non-vanishing
continuous
tangent vector field
on even-dimensional n -
spheres



Colloid-LC Composite Material in Flow

Multidimensional parameter space



Particles:

- Volume fraction
- Anchoring type: normal, planar
- Anchoring strength: strong, weak

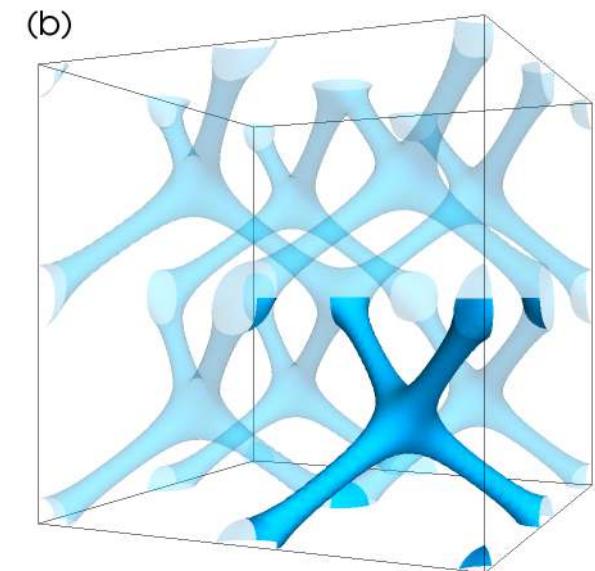
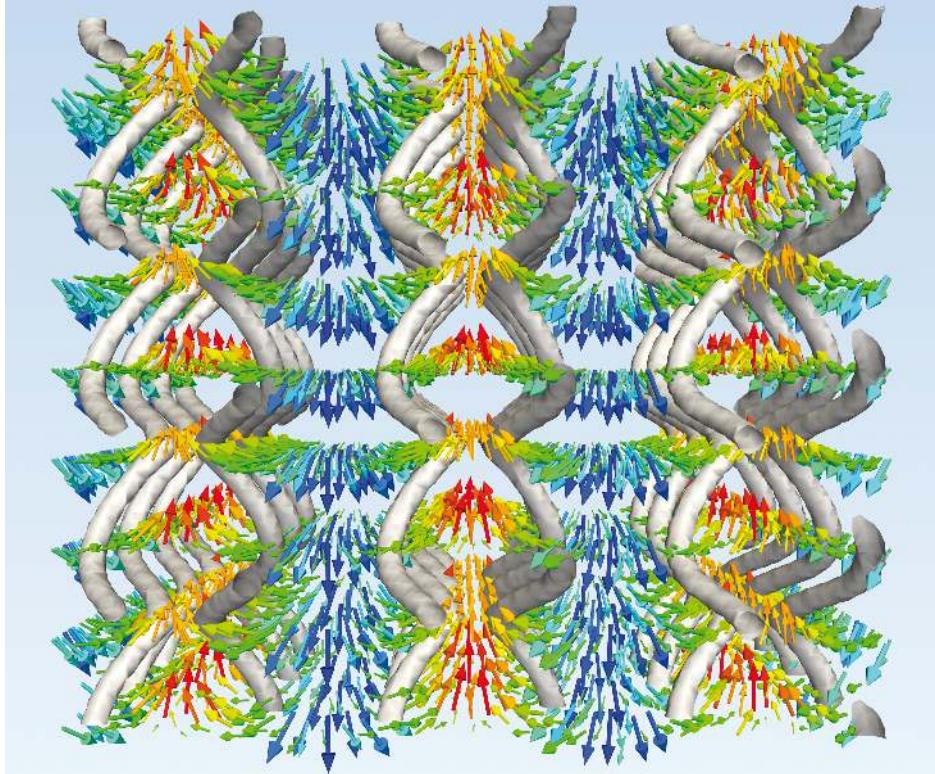
Walls:

- Anchoring type: normal, planar, hybrid
- Anchoring strength: strong, weak



Summary

- Thermodynamically consistent Landau models for complex fluids
- Mesoscopic lattice-Boltzmann simulation method
- Microfluidic flow of LC composite materials



*Flow around
disclination or defect
lines in a cholesteric
Blue Phase II (above)*



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University of Strathclyde, Glasgow

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(Department of Physics)

