

Lec 3 08 / 02 / 22 (Tue):

①

• Linear Schrödinger eqn:

$$\begin{cases} i\partial_t u = \Delta u \\ u|_{t=0} = u_0 \end{cases}$$

$$\Rightarrow u(t) = S(t) u_0$$

$$S(t) = e^{-it\Delta}$$

$$\widehat{S(t)f}(\xi) = e^{it|\xi|^2} \widehat{f}(\xi)$$

↑ smoothing (in a sense different from the heat semigroup $P(t) = e^{t\Delta}$)

① Unitary in H^s : $\| S(t)f \|_{H^s} = \| f \|_{H^s}$

② dispersive estimate:

$$\| S(t)f \|_{L_x^\infty} \lesssim \frac{1}{|t|^{d/2}} \| f \|_{L_x^1}, \quad t \neq 0$$

← "smoothing" (removes high peaks.)

Only on \mathbb{R}^d

③ Strichartz estimate:

②

We say that (q, r) is admissible if

$$\boxed{\frac{2}{q} + \frac{d}{r} = \frac{d}{2}}, \quad (q, r, d) \neq (2, \infty, 2)$$
$$2 \leq q, r \leq \infty$$

only
on \mathbb{R}^d

(i) (homogeneous), (q, r) admis

$$\| S(t) f \|_{L_t^q L_x^r} \lesssim \| f \|_{L_x^2}$$

$$\boxed{\begin{aligned} d=1: & r \leq \infty \\ d=2: & r < \infty \\ d \geq 3: & r \leq \frac{2d}{d-2} \end{aligned}}$$

(ii) (dual), (q, r) , admis

$$\| \int_{\mathbb{R}} S(-t) F(t) dt \|_{L_x^2} \lesssim \| F \|_{L_t^{q'} L_x^{r'}} \quad \begin{aligned} \frac{1}{q} + \frac{1}{q'} &= 1 \\ \frac{1}{r} + \frac{1}{r'} &= 1 \end{aligned}$$

(iii) (nonhomogeneous / retarded)

$$\| \int_0^t S(t-t') F(t') dt' \|_{L_t^q L_x^r} \lesssim \| F \|_{L_t^{q'} L_x^{\tilde{r}'}}$$

$(q, r), (q', \tilde{r}'), \text{ admis}$

(3)

• For the proof (non endpoint case), see the NLS course note
Tao's book

• $f \in L^2_x \Rightarrow S(t)f \in L^q_t L^r_x$

In particular, $S(t)f \in L^r_x$, a.s. $t \in \mathbb{R}$

• Only on $\mathbb{R}^d \Leftrightarrow$ unitarity & dispersive estimate
& H-L-S inequality (in time)

On \mathbb{T}^d , some Strichartz estimates hold.

Bourgain '93, Bourgain-Demeter '15.

• Heat eqn: Schauder estimate.

$$P(t) = e^{t\Delta}$$

$$\|P(t)f\|_{W^{s,q}} \lesssim t^{-\frac{d}{2}(\frac{1}{p}-\frac{1}{q})-\frac{s}{2}} \|f\|_{L^p_x}$$

$t > 0$

• wave eqn: Strichartz estimates.

Pf of (ii): $\Psi(t) = \int_0^t S(t-t') \Phi dW(t')$

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$$\Phi \in HS(L^2; H^s),$$

$$\Phi W = \sum_{n \in \mathbb{N}} \beta_n(t) \Phi(e_n), \quad \{e_n\}_{n \in \mathbb{N}} = \text{O.N.B of } L^2(\mathbb{R}^d)$$

$$\| \langle \nabla \rangle^s \Psi(t, x) \|_{L^2(\Omega)}$$

$$\stackrel{\text{indep}}{\underset{\text{Ho isometry}}{=}} \left\| \left(\int_0^t |S(t-t') \langle \nabla \rangle^s \Phi(e_n)(x)|^2 dt' \right)^{1/2} \right\|_{\ell_n^2}$$

$$t' \rightarrow \tau = t - t'$$

ch. of var
=

$$\| S(\tau) \langle \nabla \rangle^s \Phi(e_n)(x) \|_{\ell_n^2 L_\tau^2(\tau_0, tJ)}$$

$\forall p \geq 2$
 \Rightarrow

$$\| \langle \nabla \rangle^s \Psi(t, x) \|_{L^p(\Omega)} \lesssim p^{1/2} \| \dots \|_{L^2(\Omega)} \sim p^{1/2} \| S(\tau) \langle \nabla \rangle^s \Phi(e_n)(x) \|_{\ell_n^2 L_\tau^2(\tau_0, tJ)}$$

(5)

$$\| \|\Phi\|_{L_T^q W_x^{s,r}} \|_{L^p(\Omega)}$$

$$q, r < \infty$$

$$r \leq \frac{2d}{d-2}$$

finite $T > 0$

Minkowski: integral ineq

$$\leq \| \|\langle \nabla \rangle^s \Phi(t, x)\|_{L^p(\Omega)} \|_{L_T^q L_x^r} \quad \text{for } p \geq \frac{\max(q, r)}{r}$$

$$\leq C_p \| \|\langle \nabla \rangle^s \Phi(e_n)\|_{L_n^2 L_T^2([0, t])} \|_{L_T^q L_x^r}$$

replace by \llcorner $[0, T]$ in t over $[0, T]$

Mink

$$\leq C_p T^{1/q} \| \|\langle \nabla \rangle^s \Phi(e_n)\|_{L_T^2([0, T]; L_x^r)} \|_{l_n^2}$$

Hölder Given $2 \leq r < \infty$, let (\tilde{q}, r) be admissible for some $\tilde{q} \geq 2$

$$\leq C_p T^\theta \| \|\langle \nabla \rangle^s \Phi(e_n)\|_{L_T^{\tilde{q}}([0, T]; L_x^r)} \|_{l_n^2}$$

$$\stackrel{\text{Strichartz}}{\lesssim} C_p T^\theta \underbrace{\| \Phi(\epsilon_n) \|_{H^s} \| \cdot \|_{l^2_n}}_{= \| \Phi \|_{HS(L^2; H^s)}} < \infty \quad (6)$$

$$\Rightarrow \Psi \in L^q_T W_x^{s,r}, \text{ a.s.}$$

↑
T < ∞

□

• Local well-posedness of SNLS with additive noise

↑
existence, unique of solns,
stability under perturbation (in initial data & noise)

↑
for "short" time

$$\begin{aligned}
 \text{(SNLS)} \quad & \begin{cases} i \partial_t u - \Delta u = |u|^{2k} u + \underbrace{\Phi \sum}_{\text{additive noise}} \\ u|_{t=0} = u_0 \in H^s \end{cases}, \quad \Phi \in HS(L^2; H^s) \quad (7)
 \end{aligned}$$

We say that u is a soln to (SNLS) if u satisfies the mild formulation (= Duhamel formulation):

$$\begin{aligned}
 u(t) = & S(t) u_0 - i \int_0^t S(t-t') |u|^{2k} u(t') dt' \\
 & - i \int_0^t S(t-t') \Phi dW(t')
 \end{aligned}$$

$S(t) = e^{-it\Delta}$
 $=: \Gamma_{u_0, \Phi}(u)$

$= \Psi(t)$

① Simple case: $s > d/2$ H^s is an algebra.

Want to show $\exists! u \in C_T H^s$ s.t. $u = \Gamma(u)$. fixed pt.

⑧

(i) $\| S(t) u_0 \|_{H^s} \stackrel{\text{unitarity}}{=} \| u_0 \|_{H^s}, \quad \forall t$

$\Rightarrow \| S(t) u_0 \|_{C_T H_x^s} = \| u_0 \|_{H^s}$

(ii) $\left\| \int_0^t S(t-t') |u|^{2k} u(t') dt' \right\|_{H_x^s}$

Mink $\leq \int_0^t \| \cancel{S(t-t')} |u|^{2k} u(t') \|_{H_x^s} dt'$

unitarity $\underbrace{|u|^{2k} u(t)}_{\|u\|^{k+1} \bar{u}^k}$

product esti

\lesssim
 $\boxed{s > \frac{d}{2}}$

$\int_0^t \| u(t') \|_{H^s}^{2k+1} dt' \leq \int_0^t \| u \|_{C_T H^s}^{2k+1} dt'$
 $\forall t \in [0, T]$

$\Rightarrow \left\| \int_0^t S(t-t') |u|^{2k} u(t') dt' \right\|_{C_T H^s}$

$\lesssim T \| u \|_{C_T H^s}^{2k+1}$

(iii) Recall that

(*)

$$\| \Psi \|_{C_T H_x^s} \| L^2(\Omega) \leq C_T \| \Phi \|_{HS(L^2; H^s)}$$

$$\Rightarrow \| \Psi \|_{C_T H_x^s} \leq C(T, \| \Phi \|_{HS(L^2; H^s)}, \omega) < \infty, \text{ a.s.}$$

From (i), (ii), and (iii), we obtain, for $0 \leq T \leq 1$

$$\| \Gamma(u) \|_{C_T H_x^s} \leq \left(\| u_0 \|_{HS} + \| \Psi \|_{C([0,1]; H_x^s)} \right)$$

$$+ \left(C_1 T \| u \|_{C_T H^s}^{2k+1} \right) \leq \frac{1}{2} R \quad \text{=: } \frac{1}{2} R_\omega$$

Similarly,

$$\| \Gamma(u) - \Gamma(v) \|_{C_T H_x^s} \leq C_2 T \left(\| u \|_{C_T H^s}^{2k} + \| v \|_{C_T H^s}^{2k} \right) \| u - v \|_{C_T H_x^s} \leq \frac{1}{2}$$

$$|u|^{2k} u - |v|^{2k} v = P_{2k}(u, \bar{u}, v, \bar{v})(u-v) + Q_{2k}(u, \bar{u}, v, \bar{v})(\bar{u}-\bar{v})$$

↑
↑

poly of degree $2k$

Let $u, v \in \overline{B_R} \subset C_T H_x^s$

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$$\textcircled{A} \quad C_1 T \|u\|_{C_T H^s}^{2k+1} \leq C_1 T R^{2k+1} \leq \frac{1}{2} R$$

$$\textcircled{B} \quad C_2 T (\|u\|_{C_T H^s}^{2k} + \|v\|_{C_T H^s}^{2k}) \leq 2 C_2 T R^{2k} \leq \frac{1}{2}$$

• Banach fixed pt thm: $X =$ complete metric space.

$T =$ contraction on X : $T: X \rightarrow X$

$$d(Tx, Ty) \leq \theta d(x, y)$$

for some $0 < \theta < 1$, $\forall x, y \in X$.

Then, $\exists!$ $x \in X$ s.t. $\underline{Tx = x} \iff$ fixed pt.

Now, set $R = R(\omega) = 2 \left(\|u_0\|_{H^s} + \|\Psi\|_{C(I_0, I; H_x^s)} \right)$.

\Rightarrow choose $T = T(R\omega) > 0$ small s.t. \textcircled{A} & \textcircled{B} hold.

\Rightarrow By the Banach fixed pt thm, $\exists! u = u^\omega \in \overline{B}_{R^\omega}$ (11)

s.t. $u = \Gamma(u)$ on $[0, T_\omega]$.

• Stability?

Take $u_{0,1}, u_{0,2} \in H^s$, $\Phi_1, \Phi_2 \in HS(L^2; H^s)$

$$\Downarrow \\ \Psi_j = \int_0^t S(t-t') \Phi_j dW(t')$$

\Rightarrow Let $u_j, j=1,2$, be the solns to (SNLS):

$$\begin{cases} i \partial_t u_j - \Delta u_j = |u_j|^{2k} u_j + \Phi_j \\ u_j|_{t=0} = u_{0,j}. \end{cases}$$

A slight modification gives

(12)

$$\|u_1 - u_2\|_{C_T H_x^s} = \|\Gamma_{u_{0,1}, \Phi_1}(u_1) - \Gamma_{u_{0,2}, \Phi_2}(u_2)\|_{C_T H_x^s}$$

$$\leq \|u_{0,1} - u_{0,2}\|_{H^s} + \|\Psi_1 - \Psi_2\|_{C_T H_x^s}$$

$$+ \frac{1}{2} \|u_1 - u_2\|_{C_T H_x^s} \quad \text{by choosing } T = T_0 \text{ small.}$$

hide in LHS.

$$\Rightarrow \|u_1 - u_2\|_{C_T H_x^s} \lesssim \|u_{0,1} - u_{0,2}\|_{H^s} + \|\Psi_1 - \Psi_2\|_{C_T H_x^s}$$

Also, we have

$$\|\|\Psi_1 - \Psi_2\|_{C([0,T]; H_x^s)}\|_{L^p(\Omega)} \lesssim p^{\frac{1}{2}} \|\Phi_1 - \Phi_2\|_{HS(L^2; H^s)}$$

\forall finite $p \geq 1$

**

$$\left(\Leftarrow \Psi_1 - \Psi_2 = \int_0^t S(t-t') (\Phi_1 - \Phi_2) dW(t') \right) \quad (13)$$

By Chebyshev's inequality,

$$\| \Psi_1 - \Psi_2 \|_{C(\Sigma_0, T; H_x^s)} \leq K \| \Phi_1 - \Phi_2 \|_{HS(L^2; H^s)}$$

outside a set of probability $< c e^{-ck^2}$.

\Rightarrow LWP when $s > \frac{d}{2}$

$$u_0 \in H^s, \quad \Phi \in HS(L^2; H^s)$$

Rmk: • same proof on \mathbb{T}^d .

(14)

- At this point, we proved uniqueness of the soln only in $\overline{B_R} \subset C_T H_x^s$, but the uniqueness in fact holds in the entire $C_T H_x^s$ (unconditional uniqueness).

Method 1: Observe that if $u \in C_T H_x^s$, then

$$\lim_{t \rightarrow 0^+} \|u(t)\|_{H^s} = \|u_0\|_{H^s} \leq \frac{1}{2} R.$$

thus, by shrinking the time, we can invoke the uniqueness in $\overline{B_R}$. But, we must ensure that we can basically keep the same local existence time (up to a constant factor)

⇐ use a continuity argument / bootstrap argument

Method 2: Apply the Gronwall ineq to the difference of two given solns.

Appendix: $\textcircled{*}$ on page 9 and $\textcircled{**}$ on page 12.

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$\textcircled{*}$ on page 9:

Claim: $\| \|\Psi\|_{C_T H_x^s} \|_{L^2(\Omega)} \leq C_T \|\Phi\|_{HS(L^2; H^s)}$

Write $\Psi(t) = \Psi(t) - \overset{=0}{\Psi(0)}$

and let $F(t) = \Psi(t) / \|\Phi\|_{HS(L^2; H^s)}$.

Then, using the computation on page 14 of Lec 2, we have

$$\begin{aligned} \|F(t_1) - F(t_2)\|_{L^p(\Omega; H^s)} &\lesssim p^{1/2} \|F(t_1) - F(t_2)\|_{L^2(\Omega; H^s)} \\ &\lesssim p^{1/2} |t_1 - t_2|^{1/2}. \end{aligned}$$

Then, by Kolmogorov's conti criterion, we have

$$P \left(\sup_{0 < t \leq T} \frac{\|F(t) - F(0)\|_{H^s}}{|t|^{\alpha/p - \varepsilon}} \geq \lambda \right) \leq \frac{C_1}{\lambda^p}$$

where $p \gg 1$ and $\alpha = \frac{p}{2} - 1$

Fubini
(a.k.a. Layer-cake thm)

$$\Rightarrow \|F\|_{L^2(\Omega; C_T H_x^s)} = 2 \int_0^\infty \lambda^1 P(\|F\|_{C_T H_x^s} \geq \lambda) d\lambda$$

$$\lesssim 1 + \int_1^\infty \lambda^{1-p} d\lambda$$

for some $p \gg 1$

$< \infty$.

$$= P \left(\frac{\|F\|_{C_T H_x^s}}{T^{\alpha/p - \varepsilon}} \geq C_T \lambda \right)$$

$$\leq P \left(\frac{\|F\|_{C_T H_x^s}}{|t|^{\alpha/p - \varepsilon}} \geq C_T \lambda \right)$$

$$\leq C_T / \lambda^p$$

$$\Rightarrow \| \|F\|_{C_T H_x^s} \|_{L^2(\Omega)} \leq C_T \| \Phi \|_{HS(L^2; H^s)}$$

i.e. (*) on page 9.

- (**) on page 17;

(17)

A similar computation on pp. 15-16 gives

$$\| \|\Psi_1 - \Psi_2\|_{C(\Sigma_0, \Gamma; H_x^s)} \|_{L^p(\Omega)} \leq C_p \|\Phi_1 - \Phi_2\|_{HS(L^2; H^s)}$$

for any finite $p \geq 1$

\Rightarrow By Chebyshev's ineq,

$$\|\Psi_1 - \Psi_2\|_{C(\Sigma_0, \Gamma; H_x^s)} \leq K \|\Phi_1 - \Phi_2\|_{HS(L^2; H^s)}$$

outside a set of probability $\lesssim 1/K^p$ (a bit weaker than p. 13 but ok.)

Q: $C_p \sim p^{1/2}$: Yes, but need to work harder.

In the following, let $\Psi = \Psi_1 - \Psi_2$
 $\Phi = \Phi_1 - \Phi_2$

- Given $k \in \mathbb{Z}_{\geq 0}$, let $\{t_{l,k} : l = 0, 1, \dots, 2^k\}$ be $2^k + 1$ equally spaced pts in $[0, 1]$, i.e. $t_{0,k} = 0$
 $t_{l,k} - t_{l-1,k} = 2^{-k}$ (18)

Then, we have (with $\Psi(0) = 0$)

$$\Psi(t) = \sum_{k=1}^{\infty} \left(\Psi(t_{l_k, k}) - \Psi(t_{l_{k-1}, k-1}) \right)$$

for some $l_k = l_k(t) \in \{0, 1, \dots, 2^k\}$.

Binary expansion of $t \in [0, 1]$: $t = \sum_{k=1}^{\infty} \frac{b_k}{2^k}$, $b_k \in \{0, 1\}$
 $\Rightarrow t_{l_k, k} = \sum_{j=1}^k \frac{b_j}{2^j}$ (binary exp of t up to order k)

$$\Rightarrow \sup_{0 \leq t \leq 1} \|\Psi(t)\|_{H^s} \leq \sum_{k=1}^{\infty} \max_{0 \leq l_k \leq 2^k} \|\Psi(t_{l_k, k}) - \Psi(t_{l_{k-1}, k-1})\|_{H^s}$$

where $|t_{l_k, k} - t_{l_{k-1}, k-1}| \leq 2^{-k}$

Fix $p \gg 1$. Then, we have

$$\| \|\Psi\|_{C([0,1]; H^s)} \|_{L^p(\Omega)} \leq \sum_{k=1}^{\infty} \left\| \max_{0 \leq l_k \leq 2^k} \|\Psi(t_{l_k, k}) - \Psi(t_{l_{k-1}, k-1})\|_{H^s} \right\|_{L^p(\Omega)}$$

replace max by l^q -sum ($q \geq p$)

$$\leq \left(\int \sum_{l_k=0}^{2^k} \|\dots\|_{H^s}^q dP(\omega) \right)^{1/q} \lesssim 2^{k/q} \max_{0 \leq l_k \leq 2^k} \|\Psi(t_{l_k, k}) - \Psi(t_{l_{k-1}, k-1})\|_{L^q(\Omega; H^s)}$$

$\lesssim 1$ for $q \geq k \Rightarrow$ choose $q(k) = p + k$.

$$\lesssim \sum_{k=1}^{\infty} \max_{0 \leq l_k \leq 2^k} 2^{-\frac{k}{2}} (p+k)^{1/2} \|\Phi\|_{HS(L^2; H^s)}$$

\uparrow
 $|t_{l_k, k} - t_{l_{k-1}, k-1}|^{1/2}$

from p.14 of Lec 2.

Use $(p+k)^{1/2} \lesssim p^{1/2} \cdot k^{1/2}$.

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$$\Rightarrow \|\Psi\|_{C_T H_x^s} \Big\|_{L^p(\mathcal{Q})}$$

$$\lesssim p^{1/2} \|\Phi\|_{HS(L^2; HS)} \underbrace{\sum_{k=1}^{\infty} 2^{-k/2} \cdot k^{1/2}}_{\lesssim 1}$$

This shows ~~(*)~~ on p. 12.

(\Rightarrow Then, apply (exponential) Chebyshev to obtain the bound on p. 13.)