

Last time, we checked the regularity of the Wick power

$$:\Psi^k: = \lim_{N \rightarrow \infty} :F_N^k: , \quad F_N = P_{\leq N} \Psi_{\text{wave}}$$

for the wave case.

• We only checked $(**)$ for fixed t

but the other conditions in the regularity lemma.

(For the time difference part, simply use the mean value theorem to get $\langle m \rangle^\sigma |h|^\sigma$.

• $:\Psi^k: \in \mathcal{H}_k \Rightarrow$ By the Wiener chaos estimate, we get

$$P \left(\|:\Psi^k:\|_{L_T^q W_x^{-\varepsilon, \infty}} > \lambda \right) \leq C \exp \left(-c \frac{\lambda^{2/k}}{T^{1+2/qk}} \right)$$

$\forall \varepsilon > 0$, finite $q \geq 1$, $T > 0$.

For $q = \infty$, we can use the Garsia-Rodemich-Rumsey inequality. (2)
 and get

$$P\left(\|:\Psi^k:\|_{L^\infty([j, j+1]; W_x^{-\varepsilon, 0}(\mathbb{T}^2))} > \lambda\right) \leq C \exp\left(-c \frac{\lambda^{2/k}}{j+1}\right)$$

see Friz-Victoir
 GKO-Tolomeo
 IMRN '21

• 2-d heat: $\Psi = \Psi_{\text{heat}} = \int_0^t e^{(t-t')(\Delta-1)} dW(t')$

We only check $(**)$ in the regularity lemma. (drop N .)

$$\begin{aligned} \mathbb{E}[|\widehat{:\Psi^k:}(t, m)|^2] &= \iint_{\mathbb{T}_x^2 \times \mathbb{T}_y^2} \mathbb{E}[:\Psi^k(t, x): : \Psi^k(t, y):] \underline{e_m(y-x)} dx dy \\ &\stackrel{||}{=} k! \left\{ \mathbb{E}[\Psi(t, x)\Psi(t, y)] \right\}^k \\ &\stackrel{||}{=} \sum_m \int_0^t e^{-2(t-t')\langle m \rangle^2} dt' \underline{e_m(x-y)} \\ &= \boxed{\frac{1 - e^{-2t\langle m \rangle^2}}{\langle m \rangle^2}} \end{aligned}$$

$$\lesssim \sum_{n_1 + \dots + n_k = n} \prod_{j=1}^k \frac{1}{\langle n_j \rangle^2} \lesssim \langle n \rangle^{-2+\varepsilon}, \quad \forall \varepsilon > 0. \quad (3)$$

as in the wave case.

Prop: $\{:\Psi_N^h:\}_{N \in \mathbb{N}}$ forms a Cauchy sequence

in $L^p(\Omega; C_t C_x^{-\varepsilon}(\mathbb{T}^2))$

$$C^{-\varepsilon} = B_{\infty, \infty}^{-\varepsilon}$$

and in $C_t C_x^{-\varepsilon}(\mathbb{T}^2)$, a.s.

• Idea of the proof of the regularity lemma:

Translation invariance (in x)

$$\oplus \quad \Rightarrow \mathbb{E}[\hat{X}(t, n) \hat{X}(t, m)] = 0 \quad \text{for } n + m \neq 0.$$

(real-valued setting)

$$\begin{aligned}
& \mathbb{E} [\hat{X}(t, m) \hat{X}(t, m)] \\
&= \iint \underbrace{\mathbb{E} [X(t, x) X(t, y)]}_{= F(t, x-y)} \underbrace{e_{-n}(x) e_{-m}(y)}_{= e_{-(n+m)}(x) e_m(x-y)} dx dy \\
&= \hat{F}(t, m) \int_{\mathbb{T}^d} e^{-i(m+m) \cdot x} dx \\
&= 0 \quad \text{for } n+m \neq 0.
\end{aligned}$$

Suppose ~~(X)~~ holds. (suppress t.) Then, for $p \gg 1$,

$$\| \| X \|_{W^{s, \infty}} \|_{L^p(\Omega)} \stackrel{\text{Sobolev}}{\lesssim} \| \| X \|_{W^{s+\varepsilon, r}} \|_{L^p} \quad r \gg 1.$$

$$\stackrel{\text{Mink}}{\leq} \| \| \langle \nu \rangle^{s+\varepsilon} X(x) \|_{L^p(\Omega)} \|_{L^r_x} \quad p \geq r$$

$$\begin{aligned}
& \text{Wienerch. esti.} \\
& \leq p^{k/2} \| \| \sum_n \langle m \rangle^{s+\varepsilon} \hat{X}(m) e_n(x) \|_{L^2(\Omega)} \|_{L^r_x} \\
& \quad = \left(\bigoplus \left(\sum_m \langle m \rangle^{2(s+\varepsilon)} \mathbb{E} [|\hat{X}(m)|^2] \right) \right)^{1/2}
\end{aligned}$$

$$\textcircled{**} \lesssim p^{h/2} \left(\sum_{n \in \mathbb{Z}^d} \langle n \rangle^{2(s+\varepsilon) - d - 2s_0} \right)^{1/2} < \infty \quad \textcircled{5}$$

for $s < s_0$ (by taking $0 < \varepsilon \ll 1$ s.t. $s + \varepsilon < s_0$).

□ (As for (ii), use Kolmogorov's conti criterion.)

2.1 LWP of SNLW on \mathbb{T}^2 .

Basic product estimates:

(i) (fractional Leibniz rule): $1 < p_j, q_j, r < \infty$, $\frac{1}{p_j} + \frac{1}{q_j} = \frac{1}{r}$, $j=1,2$

$$0 \leq s \leq 1 \quad \|\langle \nabla \rangle^s (fg)\|_{L^r(\mathbb{T}^d)} \lesssim \left(\|f\|_{L^{p_1}} \|\langle \nabla \rangle^s g\|_{L^{q_1}} + \|\langle \nabla \rangle^s f\|_{L^{p_2}} \|g\|_{L^{q_2}} \right).$$

(ii) $1 < p, q, r < \infty$

$$\frac{1}{p} + \frac{1}{q} \leq \frac{1}{r} + \frac{s}{d}$$

$s > 0$

$$\|\langle \nabla \rangle^{-s} (fg)\|_{L^r(\mathbb{T}^d)} \lesssim \|\langle \nabla \rangle^{-s} f\|_{L^p} \|\langle \nabla \rangle^s g\|_{L^q}$$

SNLW:
$$\begin{cases} (\partial_t^2 + 1 - \Delta) u + u^k = \Xi & \text{on } \mathbb{T}^2. \\ (u, \partial_t u)|_{t=0} = (u_0, u_1) \end{cases}$$

(7)

truncated noise $\Xi_N = P_{\leq N} \Xi$.

Consider

$$(\partial_t^2 + 1 - \Delta) u_N + u_N^k = \Xi_N$$

↖ no truncation here.

• Duhamel

$$\Rightarrow u_N = \Psi_N + \mathcal{V}_N$$

$$\Rightarrow u_N^k = \sum_{j=0}^k \binom{k}{j} \underline{\Psi_N^j} \mathcal{V}^{k-j}$$

↓
No limit for $j \geq 2$.
need to renormalize this.

Replace $\underline{\Psi_N^j}$ by $H_j(\Psi_N; \sigma_N)$
↳ $+ \log N$.

$$\underline{H_k(x+y; \sigma)} = \sum_{j=0}^k \binom{k}{j} \underline{H_j(x; \sigma)} y^{k-j}$$

(8)

- Define the Wick power $:U_N^k:$

$$:U_N^k: = \sum_{j=0}^k \binom{k}{j} \underbrace{: \Psi_N^j :}_{\downarrow} \underbrace{V_N^{k-j}}_{\leftarrow}$$

limit exists a.s.

postulate that
the remainder is
smoother

- defined only for U_N of the form

$$U_N = \Psi_N + V_N$$

- V_N satisfies

$$(\partial_t^2 + 1 - \Delta) V_N + \sum_{j=0}^k \binom{k}{j} \underbrace{: \Psi_N^j :}_{\downarrow} V_N^{k-j} = 0.$$

$$: \Psi^j :$$

By taking $N \rightarrow \infty$, we arrive at

(SNLW_r)
$$(\partial_t^2 + 1 - \Delta) v + \sum_{j=0}^k \binom{k}{j} : \Psi^j : v^{k-j} = 0.$$

Then, the "solution" u to the renormalized SNLW:

(*)
$$(\partial_t^2 + 1 - \Delta) u + : u^k : = \xi$$

(**) is given by
$$u = \Psi + v$$

• Here, the eqn (*) for u is just a formal expression.
 For example, when $k=3$, we have

$$(\partial_t^2 + 1 - \Delta) u + u^3 - 3 \cdot \underbrace{\infty}_{= \lim_{N \rightarrow \infty} \sigma_N} \cdot u = \xi$$

By (*), we really mean (SNLW_r) with the first order decomposition (**).