

First order decomposition: $u = \Psi + v$

\uparrow
smoother

- Da Prato - Debussche '02.
- McKean '95, Bourgain '96.

Consider

$$\begin{cases} (\text{SNLW}_V') \quad | \quad (\partial_t^2 + 1 - \Delta) v + \sum_{j=0}^k \binom{k}{j} \underline{\Sigma_j} v^{k-j} = 0 \\ (\underline{\Sigma}, \partial_t v)|_{t=0} = (\underline{u}_0, \underline{u}_1), \quad \underline{\Sigma}_0 = 1. \end{cases}$$

Enhanced data set: $(\underline{u}_0, \underline{u}_1, \underline{\Sigma}_1, \dots, \underline{\Sigma}_k)$

$$\underline{\Sigma} \in \mathcal{M}^S \times \left(L^\infty([0, 1]; W^{-\varepsilon, \infty}(\mathbb{T}^2))^{\otimes k} \right)$$

$$\|\Gamma\|_{X^S} = \|(u_0, u_1)\|_{X^S} + \sum_{j=1}^k \|\Gamma_j\|_{L^\infty([0, T]; W_x^{-\varepsilon, \infty})} \quad (2)$$

Prop: $\exists \varepsilon_k > 0$ small s.t. for $0 < \varepsilon < \varepsilon_k$,

$(SNLW'_p)$ is locally well-posed:

- existence, uniqueness,
- $\Gamma \mapsto v$ is continuous.

$$C([0, T]; H^{1-\varepsilon})$$

$$T = T(\|\Gamma\|_{X^S}) > 0$$

• Case 1: $\varepsilon = 1 - \varepsilon$.

Duhamel formulation

$$v(t) = P v(t) = \partial_t S(t) u_0 + S(t) u_1 - \sum_{j=0}^k \binom{k}{j} \int_0^t S(t-t') \left(\Gamma_j v^{t-j}(t') \right) dt'$$

$$S(t) = \frac{\lim_{\tau \rightarrow t} \langle \nabla \rangle}{\langle \nabla \rangle}$$

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$$\underline{j=0}: \quad \left\| \int_0^+ S(t-t') v^k(t') dt' \right\|_{C_T H_x^{1-\varepsilon}}$$

$$\lesssim T \| v^k \|_{C_T H_x^{-\varepsilon}}$$

$$\frac{s}{d} \geq \frac{1}{p} - \frac{1}{q}$$

 $T \leq 1$

$$\stackrel{\text{Sobolev}}{\lesssim} T \| v^k \|_{C_T L_x^{2/(1+\varepsilon)}}$$

$$\frac{\varepsilon}{2} = \frac{1+\varepsilon}{2} - \frac{1}{2}$$

$$= T \| v \|_{C_T L_x^{2k/(1+\varepsilon)}}^k$$

$$\stackrel{\text{Sobolev}}{\lesssim} T \| v \|_{C_T H_x^{1-\varepsilon}}^k$$

$$\frac{1-\varepsilon}{2} \geq \frac{1}{2} - \frac{1+\varepsilon}{2k}$$

$$\left(\Leftrightarrow \varepsilon \leq \frac{1}{k-1} \right)$$

$$\underline{1 \leq j < k}: \quad \left\| \int_0^t S(t-t') (\sum_i v^{k-i}) dt' \right\|_{C_T H_x^{1-\varepsilon}}$$

$$\lesssim T \left\| \sum_i v^{k-i} \right\|_{C_T H_x^{-\varepsilon}}$$

↑
only a distribution (in x)

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$$= T \left\| \langle \nabla \rangle^{\varepsilon} (\sum_j V^{k-j}) \right\|_{C_T L_x^2}$$

$$\lesssim T \left\| \langle \nabla \rangle^{\varepsilon} \sum_j \right\|_{L_T^\infty L_x^{2/\varepsilon}} \left\| \langle \nabla \rangle^{\varepsilon} (V^{k-j}) \right\|_{L_T^\infty L_x^2}$$

frac. Leib. rule

$$\lesssim T \left\| \sum_j \right\|_{X^{1-\varepsilon}} \left\| \langle \nabla \rangle^{\varepsilon} V \right\|_{L_T^\infty L_x^{2(k-j)}}^{k-j} \quad \frac{1}{p} + \frac{1}{2} \leq \frac{1}{2} + \frac{\varepsilon}{2}$$

Sobolev

$$\lesssim T \left\| \sum_j \right\|_{X^{1-\varepsilon}} \|V\|_{C_T H_x^{1-\varepsilon}}^{k-j} \quad \frac{1-2\varepsilon}{2} \geq \frac{1}{2} - \frac{1}{2(k-j)}$$

 $\tilde{J} = h$:

$$\left\| \int_0^t S(t-t') \sum_k F_k(t') dt' \right\|_{C_T H_x^{1-\varepsilon}}$$

$$\lesssim T \left\| \sum_k \right\|_{C_T \tilde{H}_x^{-\varepsilon}}$$

$$\lesssim T \left\| \sum_k \right\|_{X^{1-\varepsilon}}$$

\Rightarrow Putting together, we get

$$\begin{aligned}\|\Gamma v\|_{C_T H_x^{1-\varepsilon}} &\leq C_1 \|(\bar{u}_0, \bar{u}_1)\|_{H^{1-\varepsilon}} \\ &+ C_2 T \underbrace{\|\mathcal{E}\|_{X^{1-\varepsilon}}}_{k^*} \left(1 + \|v\|_{C_T H_x^{1-\varepsilon}}\right)^{k-1} \\ &+ C_3 T \|v\|_{C_T H_x^{1-\varepsilon}}^k\end{aligned}$$

A similar estimate holds for the difference $\Gamma(v_1) - \Gamma(v_2)$.

$\Rightarrow \Gamma$ is a contraction on a ball of radius

$$\sim \|(\bar{u}_0, \bar{u}_1)\|_{H^{1-\varepsilon}} + \|\mathcal{E}_k\|_{C([0,1]; W_x^{-\varepsilon, \omega})}$$

by choosing $T = T(\|\mathcal{E}\|_{X^{1-\varepsilon}}) > 0$ sufficiently small.

Summary :

Known : 2-d NLW is ill-posed in negative Sobolev spaces
 (christ - Colliander - Tao, Oh - Okamoto - Tzvetkov,
 '03
 Forlano - Okamoto DPDE '20)

\Rightarrow 2-d SNLW:

$$(u_0, u_1, \vec{z}) \rightarrow u \text{ ill-posed.}$$

Idea: Decompose the ill-defined soln map.

① Use stochastic analysis to construct an enhanced data set only in this step.

② use deterministic analysis to prove LWP for the remainder term $v = u - \Psi$.

continuous map

$$(u_0, u_1, \vec{z}) \xrightarrow{\textcircled{1}} \Xi = (u_0, u_1, \Psi, : \Psi^2 : , \dots , : \Psi^k :) \xrightarrow{\textcircled{2}} v \rightarrow u = \Psi + v$$

↑
NOT conti

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Thm: $\exists \varepsilon_0 > 0$ small s.t. for $0 < \varepsilon < \varepsilon_0$,

the renormalized SNLW on \mathbb{T}^2 is locally well-posed in $H^{1-\varepsilon}(\mathbb{T}^2)$.

Moreover, the soln u_N to the truncated renormalized

$$\text{SNLW : } (\partial_t^2 + 1 - \Delta) u_N + :u_N^\ell: = \mathcal{F}_N$$

converges to $u = \Psi + v$. in $C([\Sigma_0, T_0]; H^{-\varepsilon}(\mathbb{T}^2))$, a.s.

- First order expansion $u = \underset{\leftarrow \text{rough}}{\Psi} + \underset{\leftarrow \text{smoother}}{v}$

- trick to solve the equation
- also gives a description

ex: in small scales, u "behaves like" Ψ .

- Without renormalization, no non-trivial limit exists.

($u_N \rightarrow 0$ or \lim soln.)

Oh-Okamoto Robert

SPA '20.

\Leftarrow Triviality

• Regulation on $\tilde{\gamma}$:

- by $P \in N$ $\tilde{\gamma}_\delta = \gamma_\delta * \tilde{\gamma}$ (conv in prob.)
 - by mollification (spatial / space-time).
- \Rightarrow same limit

Q: Rougher initial data?

$$\text{NLW: } \partial_t^2 u - \Delta u + u^k = 0 \quad \text{on } \mathbb{R}^d.$$

• Scaling invariance

\Rightarrow Scaling critical regularity

$$s_{\text{scaling}} = \frac{d}{2} - \frac{2}{k-1}.$$

• Lorentz invariance \Rightarrow $s_{\text{conf}} = \frac{d+1}{4} - \frac{1}{k-1}$
(Conformal)

$$s_{\text{crit}} = \max(s_{\text{scaling}}, s_{\text{conf}}, 0)$$

$$\underset{d=2}{=} \max\left(1 - \frac{2}{k-1}, \frac{3}{4} - \frac{1}{k-1}, 0\right).$$

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Thm: GKO TAMS '18

$$(i) k=2,3: s > s_{\text{crit}} \quad (ii) k \geq 4: s \geq s_{\text{crit}}.$$

The renormalized SNLW on \mathbb{T}^2 is locally well-posed in $H^s(\mathbb{T}^2)$.

Schwarz estimates: $0 \leq s \leq 1$

for initial data

(q, r) s -admissible

$$\text{if } 1 \leq \tilde{q} \leq z \leq q \leq \infty$$

(\tilde{q}, \tilde{r}) dual s -admissible

$$1 \leq \tilde{r} \leq z \leq r \leq \infty$$

$$\text{scaling: } \frac{1}{q} + \frac{d}{r} = \frac{d}{2} - s = \frac{1}{\tilde{q}} + \frac{d}{\tilde{r}} - 2$$

$$(q, r, d) \neq (2, \infty, 3)$$

$$\text{admissibility: } \frac{1}{q} + \frac{d-1}{2r} \leq \frac{d-1}{4}, \quad \frac{1}{\tilde{q}} + \frac{d-1}{2\tilde{r}} \leq \frac{d-1}{4}$$

$$(\tilde{q}, \tilde{r}, d) \neq (2, 1, 3)$$

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(\tilde{g}, \tilde{r}) dual s -admissible

$\Leftrightarrow (\tilde{g}', \tilde{r}')$ is $(1-s)$ -admissible

Then, $\begin{cases} (\partial_t^2 + 1 - \Delta) u = F \\ (u, \partial_t u)|_{t=0} = (u_0, u_1) \end{cases}$ on \mathbb{T}^d .

$$\Rightarrow \| (u, \partial_t u) \|_{L_T^{10} H_x^s} + \| u \|_{L_T^q L_x^r} \quad 0 \leq T \leq 1$$

$$\lesssim \| (u_0, u_1) \|_{H_x^s} + \| F \|_{L_T^{\frac{q}{2}} L_x^{\frac{r}{2}}}$$

Also, $(LHS) \lesssim \| (u_0, u_1) \|_{H_x^s} + \| F \|_{L_T^1 H_x^{s-1}}$

- Strichartz on \mathbb{R}^d : Visan's Oberwolfach note.
and finite speed of propagation

- We only consider $k=3$. $\Rightarrow \text{Suit} = 1/4$.

(general case : GKO, Oh-Thomann, Toulouse'20.

$$(q, r) = \left(\frac{12}{1+4\delta}, \frac{3}{1-2\delta} \right), \quad (\frac{1}{4} + \delta) - \text{admis.}$$

$$(\tilde{q}, \tilde{r}) = \left(\frac{12}{q+4\delta}, \frac{3}{3-2\delta} \right), \quad \text{dual } (\frac{1}{4} + \delta) - \text{admis.}$$

Nonlinearity: $\sum_j v^{3-j} \cdot \text{in } L_T^{\tilde{q}} L_x^{\tilde{r}} + L_T^1 H_x^{s-1}$

$$\|v^3\|_{L_T^{\tilde{q}} L_x^{\tilde{r}}} = \|v\|_{L_T^{\frac{36}{q+4\delta}} L_x^{\frac{9}{3-2\delta}}}^3 \lesssim T^\theta \|v\|_{L_T^q L_x^r}^3$$

$$\frac{36}{q+4\delta} < q = \frac{12}{1+4\delta}$$

$$\frac{9}{3-2\delta} < r = \frac{3}{1-2\delta}$$

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deterministic theory.

LWP of cubic NLW
in $H^s(\mathbb{T}^2)$, $s > \frac{1}{4}$.

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$$\cdot \|\Sigma_3\|_{L_T^1 H_x^{s-1}} \leq T \|\Sigma_3\|_{L_T^\infty W_x^{-\varepsilon, \infty}} \quad \text{if } s-1 \leq -\varepsilon.$$

$$\leq T \|\Sigma\|_{H^s}.$$

$$s = \frac{1}{4} + \delta$$

$$\cdot \|\Sigma_j v^{3-\frac{j}{d}}\|_{L_T^1 H_x^{s-1}} \quad s-1 = -\frac{3}{4} + \delta$$

$$\stackrel{\text{Sobolev}}{\leq} \|\langle \nabla \rangle^{-\delta} (\Sigma_j v^{3-\frac{j}{d}})\|_{L_T^1 L_x^{\frac{8}{7-8\delta}}} \quad \frac{\frac{3}{4} - 2\delta}{2} = \frac{7-8\delta}{8} - \frac{1}{2}$$

$$\lesssim \|\langle \nabla \rangle^{-\delta} \Sigma_j\|_{L_T^\infty L_x^\infty} \|\langle \nabla \rangle^{\delta} v^{3-\frac{j}{d}}\|_{L_T^1 L_x^{\frac{8}{7-8\delta}}} \quad \frac{1}{p} + \frac{1}{q} \leq \frac{1}{r} + \frac{s}{d}$$

replace by L_x^∞

$$\frac{\delta}{2} + \frac{7-8\delta}{2} \leq \frac{7-8\delta}{8} + \frac{\delta}{2}$$

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$$\hat{j}=2: \lesssim T \|v\|_{L_T^\infty H_x^{\frac{d}{2}}} \leq T \|v\|_{L_T^\infty H_x^s}$$

$$\hat{j}=1: \lesssim T \|\langle \nabla \rangle^{\delta} v\|_{L_T^\infty L_x^{\frac{16}{7-8\delta}}} \stackrel{\text{Sobolev}}{\lesssim} T \|v\|_{L_T^\infty H^{\frac{1+16\delta}{8}}}^2$$

frac Leib

$$\frac{1+\frac{8\delta}{8}}{8}/2 = \frac{1}{2} - \frac{7-8\delta}{16} \quad \text{and add } \delta \text{ from } \langle \nabla \rangle^{\delta}$$

$$\text{Let } \|\vec{u}\|_{Y^s(\Gamma)} = \|(\psi, \partial_\nu \psi)\|_{L_T^{\frac{10}{3}} H_x^s} + \|u\|_{L_T^q L_x^r}.$$

$$\begin{aligned} \Rightarrow \|\vec{\Gamma}(v)\|_{Y^s(\Gamma)} &= \|\left(\Gamma(v), \partial_\nu \Gamma(v)\right)\|_{Y^s(\Gamma)} \\ &\leq c_1 \|(\psi_0, u_0)\|_{H^s} + c_2 \left(\|v\|_{Y^s(\Gamma)} + \|\Xi\|_{\Sigma^{-s}} \right)^3 \end{aligned}$$

$$\text{where } \|\Xi\|_{\Sigma^{-s}} = \sum_{j=1}^3 \|\Xi_j\|_{L^{\infty}([0,1]; W_x^{-s, \infty})}$$

A similar estimate holds for the difference.