

II WHITE FORCED NSE

$$f = \frac{d}{dt} z(t, x)$$

$$z(t, x) = \sum_{n \in \mathbb{Z}_0^2} b_n \beta_n(t) e_n(x)$$

← ind. redl BM

div free

STOCHASTIC CONVOLUTION

$$\Psi(t) = \int_0^t e^{(t-t')L} d z(t'), \quad L = \mathbb{P} \Delta$$

• regularity: $z = \phi W$ ← div free version

$$\|\phi\|_{HS(L^2; H^s)} = \|\phi(e_n)\|_{\ell_n^2 H^s} = \left(\sum_n |n|^{2s} b_n^2 \right)^{1/2} =: \sqrt{B_s}$$

$$\phi e_n = b_n e_n$$

If $\phi \in HS(L^2; H^s)$, then $z \in C_t W_x^{s+1-\epsilon, \infty}$, $\forall \epsilon > 0$

$z \in L_T^q W_x^{s+1, r}$, $T < \infty$ a.s.
 ↑
 no loss $q, r < \infty$

WELL-POSEDNESS OF (SNSE)

$$\partial_t u = Lu - B(u) + \partial_t z \quad | \quad du = (Lu - B(u)) dt + dz$$

• Da Prato - Debussche trick ('03)

← but dates back further

$$u = v + \Psi$$

$$\text{Recall } \partial_t \Psi = L\Psi + \partial_t z$$

$$v = u - \Psi$$

$$\Rightarrow \partial_t v = Lv - \underbrace{B(v, v) + B(v, \Psi) + B(\Psi, v) + B(\Psi, \Psi)}_{\text{"forcing term"}} \quad (\text{SNSE}')$$

• A priori estimate

(SNSE') $\times v$

$$\bullet \langle B(\Psi, v), v \rangle = 0$$

$$\bullet \langle B(v, \Psi), v \rangle \leq c_1 \|v\|_{L^4}^2 \|\nabla \Psi\|_{L^2} \leq c_2 \|v\|_{L^2} \|v\|_{H^1} \|\nabla \Psi\|_{L^2}$$

$$\leq \frac{1}{2} \|v\|_{H^1}^2 + c_3 \|v\|_{L^2}^2 \|\nabla \Psi\|_{L^2}^2$$

$$\bullet \langle B(\Psi, \Psi), v \rangle \lesssim \|v\|_{L^2}^2 + \|\langle \nabla \rangle \Psi\|_{L^4}^4$$

$$\Rightarrow \partial_t \|v\|_{L^2}^2 \leq c_1(\Psi(t)) + c_2(\Psi(t)) \|v\|_{L^2}^2$$

Gronwall's

$$\Rightarrow \sup_{t \in [0, T]} \|v(t)\|_{L^2} \leq e(\nu_0, \Psi, T)$$

Also, (SNSE') $\times v$
 $\Rightarrow \|v\|_{L_T^2 H_x^1} \leq e(\nu_0, \Psi, T)$

$$\bullet \|\mathcal{B}(v, \Psi)\|_{L_T^2 H_x^{-1}}^2 \lesssim \|v\|_{L_T^\infty L_x^2}^2 \|\Psi\|_{L_T^2 L_x^\infty}^2$$

$$\langle \mathcal{B}(v, \Psi), w \rangle = - \langle \mathcal{B}(v, w), \Psi \rangle \sim \int_{L^2} v \cdot \nabla w \cdot \Psi_{L^2 L^\infty}$$

$$\bullet \|\mathcal{B}(\Psi, v)\| \lesssim \|v\|_{L_T^\infty L_x^2} \|\Psi\|_{L_T^2 L_x^\infty}$$

$$\bullet \|\mathcal{B}(\Psi, \Psi)\|_{L_T^2 H_x^{-1}} \lesssim \|\Psi\|_{L_T^4 H_x^{1/2}}^2$$

$$\Rightarrow \|\partial_t v\|_{L_T^2 H_x^{-1}} \lesssim \|v\|_{L_T^2 H_x^1} + \|v\|_{L_T^\infty L_x^2} \|v\|_{L_T^2 H_x^1} + \|\Psi\|_{L_T^2 L_x^\infty} \|v\|_{L_T^\infty L_x^2} + \|\Psi\|_{L_T^4 H_x^{1/2}}^2$$

$$\Rightarrow \|\lambda u\|_{L_T^2 H_x^{-1}} \leq C(\mu_0, \Psi, \tau)$$

GALERKIN APPROXIMATION

$$\begin{cases} \partial_t v_N = L v_N - P_N \mathcal{B}(v_N + \Psi_N) & , \quad \Psi_N = P_N \Psi \\ v_N|_{t=0} = P_N \mu_0 \end{cases}$$

→ POINT: same a priori estimate (unif. in N)

$$\Psi_N \rightarrow \Psi \quad \text{a.s.}$$

$$(\Rightarrow \mathcal{B}(v_N + \Psi_N) \rightarrow \mathcal{B}(v + \Psi) \text{ in } L_T^1 H_x^{-1})$$

result follows as before

⇒ well-posed on $[0, \tau]$, $\forall \tau > 0$

⇒ GWP in $L^2(\mathbb{T})$ if $B_0 < \infty$ (i.e. $\phi \in \text{HS}(L^2; L^2)$)

This argument shows a control on $\|u(t)\|_{L^2} \leq \|v(t)\|_{L^2} + \|\Psi(t)\|_{L^2}$.

But we can use Itô's Lemma, as well

$$dx^{(i)} = \sum_{j=1}^m f_{ij} dB_j + g_i dt$$

$$B = \begin{pmatrix} B_1 \\ \vdots \\ B_m \end{pmatrix}_{\text{ind.}}, \quad \vec{x} = \begin{pmatrix} x^{(1)} \\ \vdots \\ x^{(n)} \end{pmatrix}, \quad f = (f_{ij})_{n \times m}, \quad g = \begin{pmatrix} g_1 \\ \vdots \\ g_n \end{pmatrix}$$

Then,

$$dF(t, \vec{x}_t) = \frac{\partial F}{\partial t}(t, \vec{x}_t) dt + \sum_{i=1}^m \frac{\partial F}{\partial x_i}(t, \vec{x}_t) dx_t^{(i)} \leftarrow \nabla_{\vec{x}} F \cdot d\vec{x}_t$$

$$+ \frac{1}{2} \sum_{i,j=1}^m \frac{\partial^2 F}{\partial x_i \partial x_j}(t, \vec{x}_t) dx^{(i)} dx^{(j)}$$

$$\frac{1}{2} \nabla^2 F(d\vec{x}_t, d\vec{x}_t)$$

$\frac{dB_j}{dt}$	$\frac{dB_j}{dt}$	$\frac{dB_j}{dt}$
$\frac{dB_j}{dt}$	$\frac{dB_j}{dt}$	$\frac{dB_j}{dt}$
$\frac{dB_j}{dt}$	$\frac{dB_j}{dt}$	$\frac{dB_j}{dt}$

$$du = V_N(u) dt + \sum_{n \in \mathbb{N}} b_n e_n d\beta_n$$

$$\Rightarrow dF(t, u(t)) = \partial_t F(t, u(t)) dt + \langle \nabla_u F, V_N(u) \rangle dt + \langle \nabla_u F, \sum_{n \in \mathbb{N}} b_n e_n d\beta_n \rangle + \frac{1}{2} \sum_{n \in \mathbb{N}} \frac{\partial^2 F}{\partial \hat{u}_n^2} b_n^2 \cdot dt$$

with $F(t, u) = \|u\|_{L^2}^2$.

Apply Itô's Lemma and take expectation on $\mathbb{N} \in \mathbb{N}$

$$\Rightarrow \frac{d}{dt} \mathbb{E} [\|u\|_{L^2}^2] = -2 \mathbb{E} (\|u\|_{H^1}^2) + B_{0,N} \leq -2 \mathbb{E} (\|u\|_{L^2}^2) + B_0$$

→ Using Gronwall \Rightarrow control on $\mathbb{E} [\|u(t)\|_{L^2}^2]$ for all fixed t . But we want to control $\mathbb{E} [\sup_{t \in [0, T]} \|u(t)\|_{L^2}^2]$.

Key Tool: Burkholder-Davis-Gundy Inequality (also called Martingale inequality)

$$\left(\mathbb{E} \left[\sup_{0 \leq t \leq T} |M(t)|^p \right] \sim \mathbb{E} \left[\langle M \rangle_T^{p/2} \right], p > 0 \right.$$

where $M(t)$ is a local martingale. quadratic variation

think of M as $\int_0^t F dB_t$

$$\Rightarrow \mathbb{E} \left[\sup_{0 \leq t \leq T} \|u(t)\|_{L^2}^2 \right] \leq C (\|u_0\|_{L^2}^2, B_0, T)$$

By applying Itô's Lemma to $F(u) = \|u\|_{H^1}^2$

$$\Rightarrow \mathbb{E} [\|u(t)\|_{H^1}^2] \leq \frac{1}{2} B_1 + e^{-2t} \mathbb{E} [\|u_0\|_{H^1}^2]$$

STATIONARY MEASURE ?

Semigroup: $T_t : C_b(H) \rightarrow C_b(H)$, $T_t^* : P(H) \rightarrow P(H)$

$$(T_t f)(v) = \int_H f(u) P_t(v, du) = \mathbb{E} [f(u(t; v))] \\ (T_t^* \mu)(A) = \int_H P_t(v, A) \mu(dv) = (\Phi_t)_* \mu$$

push-forward Φ_t solution map

If $\mathcal{L}(\mu_0) = \mu$, then $\mathcal{L}(u(t; \mu_0)) = (T_t^* \mu)$

Theorem: \exists invariant measure.

proof: For simplicity, assume $B_1 < \infty$. The result also holds for $B_1 = \infty$.

$u(t)$ = solution to (SNSE) with $u_0 = 0$

$\mu_t = \mathcal{L}(u(t))$

$\bar{\mu}_t = \frac{1}{t} \int_0^t \mu_{t'} dt'$

$B_{H^1}(r)$ = ball of radius r in H^1 (centered at 0)

$$\mu_t (H^1 B_{H^1}(r)) = P(\|u(t)\|_{H^1} > r)$$

$$\text{Chebyshev} \leq \frac{C}{r^2}$$

$$< \varepsilon$$

by choosing $r \gg 1$ for a given $\varepsilon > 0$, uniformly in $t \geq 0$

$$\Rightarrow \bar{\mu}_\varepsilon (B_{H^1}(r))^c < \varepsilon$$

$$\Rightarrow \{\bar{\mu}_t\} \text{ is tight } \Rightarrow \dots \text{ rest as before. } \square$$

"UNIVERSALITY" OF WHITE-NOISE FORCES

consider random kick forces

$$\eta_\varepsilon(t) = \sqrt{\varepsilon} \sum_{k \in \mathbb{Z}_0} \eta_k^\omega \delta(t - \varepsilon k) \quad \text{where}$$

$$\eta_k = \sum_{n \in \mathbb{Z}_0^d} b_n g_{kn} e_n$$

Compare the dynamics with (SINSE) where $f = \frac{d}{dt} \bar{z}$, $\bar{z} = \sum_{n \in \mathbb{Z}_0^d} b_n \beta_n e_n$.

claim: " μ^ε converges to μ in law"

\Leftarrow Donsker's Theorem

$\{X_n\}_n$ i.i.d. mean 0, invariance σ^2

$$Z_n(t; \omega) = \frac{1}{\sigma \sqrt{n}} \int_{[nt]}^t (\omega) + (nt - [nt]) \frac{1}{\sigma \sqrt{n}} X_{[nt]+1}(\omega)$$

$$\downarrow \sum_{j=1}^{[nt]} X_j$$

BM in law

- Remark:
- Can also consider continuous in time forcing but not Gaussian, and still obtain "weak universality" and convergence to white forcing.
 - Skorokhod's theorem can be used to upgrade convergence in law to a.s. convergence.