

Goal: Disprove the freq loc. dual estimate:

$$\textcircled{*} \quad \left\| \int e^{it\sqrt{\Delta}} F(t) dt \right\|_{L_x^2} \lesssim \|F\|_{L_t^2 L_x^1}$$

β_t = Brownian motion

$\beta_t - \beta_s$ = mean 0 Gaussian r.v.
with Var = $t-s$

$$\mathbb{E} \left[e^{\pm 2\pi i (\beta_t - \beta_s) \xi} \right] = e^{-2\pi |t-s| \xi^2}, \quad \forall \xi \in \mathbb{R}$$

$$\text{Let } F(t, x) = \gamma\left(\frac{t}{T}\right) \psi(x - t \hat{e}_3 - \beta_t \hat{e}_1)$$

↑
non-neg, smooth cutoff
supp on $[-1, 1]$

$\hat{\psi}(\xi)$ supported
on $\{|\xi| \sim 1\}$

Suppose $\textcircled{*}$ holds

$$\text{With } \hat{F}(t, \xi) = \gamma\left(\frac{t}{T}\right) e^{-2\pi i t \xi_3 - 2\pi i \beta_t \xi_1} \hat{\psi}(\xi)$$

$$\textcircled{*} \Rightarrow \left\| \int e^{2\pi i t |\vec{\beta}|} e^{-2\pi i t \vec{\beta}_3} e^{-2\pi i C \beta t} \gamma\left(\frac{t}{T}\right) dt \cdot \hat{\psi}(\vec{\beta}) \right\|_{L^2_{\vec{\beta}}} \quad \textcircled{2}$$

$\lesssim \|F\|_{L_t^2 L_x^1} \sim \sqrt{T}.$

$\textcircled{**}$

$$\int_{|\vec{\beta}| \sim 1} \left| \int_t \dots dt \right|^2 d\vec{\beta} \lesssim T$$

Let $A = \{ \vec{\beta} \in \mathbb{R}^3 : |\vec{\beta}_3| \sim 1, |\vec{\beta}_2| \ll |\vec{\beta}_1| \ll 1 \}$.

\Rightarrow We have the following Taylor approx:

$$|\vec{\beta}| - \vec{\beta}_3 \sim \vec{\beta}_1^2, \quad \forall \vec{\beta} \in A.$$

$$\left(|\vec{\beta}| = \sqrt{\vec{\beta}_3^2 + \vec{\beta}_1^2 + \vec{\beta}_2^2} = \vec{\beta}_3 + \frac{1}{2|\vec{\beta}_3|} (\vec{\beta}_1^2 + \vec{\beta}_2^2) + \text{error} \right)$$

- Restrict the domain to A and expand the square.

(3)

(LHS) of $\star\star$

$$\sim \int_A \iint_{t+t'} e^{2\pi i(t-t')(|\xi| - |\xi_3|)} \frac{e^{-2\pi i(\beta_t - \beta_{t'})|\xi_1|}}{\gamma\left(\frac{t}{T}\right)\gamma\left(\frac{t'}{T}\right)} dt dt' d\xi$$

Take \mathbb{E}

$$\int_A \iint_{t+t'} e^{2\pi i(t-t')(|\xi| - |\xi_3|)} \bar{e}^{-2\pi|r-t'||\xi|^2} \gamma\left(\frac{t}{T}\right)\gamma\left(\frac{t'}{T}\right) dt dt' d\xi \lesssim T.$$

Let $r = t - t'$.

$$\int_A \int e^{2\pi i r (|\xi| - |\xi_3|)} e^{-2\pi|r||\xi|^2} \gamma * \gamma\left(\frac{r}{T}\right) dr d\xi \lesssim 1.$$

Let $T \rightarrow \infty$

$$\int_A \int e^{2\pi i r \underbrace{(|\xi| - |\xi_3|)}_{\propto}} \bar{e}^{-2\pi|r||\xi|^2} dr d\xi \lesssim 1$$

$$\int_A \frac{1}{|\xi_1|^2} d\xi_1 d\xi_2 d\xi_3 \sim \infty$$

\uparrow
 $|\xi_2| \ll |\xi_1| \ll 1 \Rightarrow \int \frac{1}{\xi_1} d\xi_2 = \infty$

$$\begin{aligned} & \left(\int e^{2\pi i rx} = \bar{e}^{-2\pi|r|y} = dr \right. \\ & \quad \left. = \frac{y}{\pi(x^2 + y^2)} \right) \Rightarrow \text{contradiction.} \end{aligned}$$

Rmk: Without the BM factor, there would be NO $e^{-2\pi i \langle H, \vec{z}_1^2 \rangle}$ in the last integral.

$$\Rightarrow \int_A \delta_0(\vec{z}_1^2) d\vec{z} \lesssim 1 \quad (\Leftarrow \text{No contradiction})$$

Recall $\|f\|_{BMO} = \sup_Q \frac{1}{|Q|} \int_Q |f - f_Q|$, $f_Q = \operatorname{ave}_Q f$.

Lemma 2.11: $\operatorname{supp} \hat{f} \subset \{|z| \sim 2^h\}$, $h \in \mathbb{Z}$ (homog LP localization)

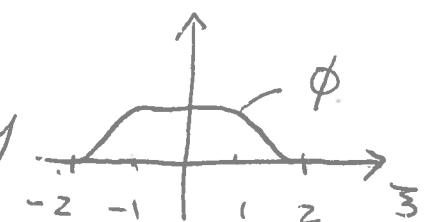
$$\Rightarrow \|f\|_{BMO} \sim \|f\|_{L^\infty}$$

Pf: By scaling, we may assume $k=0$.

• Fix x .

$$j \gg 1. \quad f(x) = P_{\leq j} f(x) = \int f(x-y) 2^{jd} \phi(2^j y) dy$$

$$\phi(z) = \begin{cases} 1, & |z| \leq 1 \\ 0, & |z| \geq 2 \end{cases}$$



We also have

(5)

$$\int f(x-y) 2^{-jd} \Phi(2^{-j}y) dy = P_{\leq -j} f(x) = 0.$$

$$\Rightarrow f(x) = \int f(x-y) \left(2^{jd} \Phi(2^jy) - 2^{-jd} \Phi(2^{-j}y) \right) dy \quad \text{proj onto } \{|y| \lesssim 2^{-j}\}.$$

Let $Q = \{y : |y| \leq 2^{\alpha}\}$. and $Q_x = x + Q$.

Note that

$$\begin{aligned} & \int f_{Q_x} (2^{jd} \Phi(2^jy) - 2^{-jd} \Phi(2^{-j}y)) dy \\ &= f_{Q_x} \left[\underbrace{\int 2^{jd} \Phi(2^jy) dy}_{= \int \Phi dy} - \underbrace{\int 2^{-jd} \Phi(2^{-j}y) dy}_{= \int \Phi dy} \right] = 0 \end{aligned}$$

$$\Rightarrow f(x) = \int (f(x-y) - f_{Q_x}) (2^{jd} \Phi(2^jy) - 2^{-jd} \Phi(2^{-j}y)) dy$$

If $y \in Q$, then $|f(x)| \lesssim 2^{d\alpha} \|f\|_{BMO}$.

If $y \notin Q$, then $|f(x)| \leq 2 \|f\|_{L^\infty}$

$$\Rightarrow |f(x)| \leq C_\alpha \|f\|_{BMO} + \frac{1}{2} \|f\|_{L^\infty}$$

\Rightarrow Take sup in x .



$$\begin{aligned} & \int |2^{jd} \Phi(2^jy)| + |2^{-jd} \Phi(2^{-j}y)| dy \\ & \underbrace{\quad |y| \geq 2^\alpha \quad}_{< \frac{1}{4}} \quad \text{by choosing } \alpha \gg 1 \quad (\text{in dep of } x) \\ & \quad \quad \quad (j \text{ is fixed}) \end{aligned}$$

(6)

Chap 3: Invariant Gibbs measure, Part 2

We'll focus on \mathbb{T}^2 .

Sec 3.1: Gibbs meas on \mathbb{T}^2 (for NLS, NLW)

$$H(u) = \frac{1}{2} \int |\nabla u|^2 \pm \frac{1}{p+1} \int |u|^{p+1}$$

$$\begin{aligned} \Rightarrow "d\mu &= Z^{-1} e^{-H(u)} du \\ &= Z^{-1} e^{\mp \frac{1}{p+1} \int |u|^{p+1}} d\mu, \end{aligned}$$

where $d\mu = Z^{-1} e^{-\frac{1}{2} \|u\|^2_{H^1}} du \leftarrow$ supported on $H^s(\mathbb{T}^2) \setminus L^2(\mathbb{T}^2)$

$$\Rightarrow \int |u|^{p+1} dx = \infty, \text{ a.s.} \quad s < 0.$$

- We need to renormalize $\int |u|^{p+1}$.

We only consider the defocusing case.

Focusing case : Brydges - Slade '96.

We can not construct Gibbs measure for the focusing cubic NLS on \mathbb{T}^2 even if we consider the Wick ordered nonlinearity (with the Wick ordered

In the following, we focus on the real-valued setting

(7)

- Under P_1 , we have $\text{indep, std Gaussian r.v.}$ (for simplicity)

$$u = \sum_{n \in \mathbb{Z}^2} \frac{g_n(\omega)}{\langle n \rangle} e^{inx}, \quad g_{-n} = \overline{g_n}, \quad \text{Var}(g_n) = 1$$

Given $N \in \mathbb{N}$, let $u_N = P_{\leq N} u$.

\Rightarrow For each $x \in \mathbb{T}^2$, $u_N(x)$ is a mean 0 \mathbb{R} -valued Gaussian r.v. with variance:

$$\sigma_N := \mathbb{E}[u_N^2(x)] = \sum_{|n| \leq N} \frac{1}{\langle n \rangle^2} \sim \log N \rightarrow \infty.$$

Similarly,

$$\mathbb{E}\left[\int_{\mathbb{T}^2} u_N^2(x) dx\right] = \sum_{|n| \leq N} \frac{1}{\langle n \rangle^2} \sim \log N$$

but defining : $u_N^2(x) := u_N^2(x) - \sigma_N$

Wick ordered monomial $\xrightarrow{\longrightarrow} = H_2(u_N(x); \sigma_N)$

\uparrow Hermite poly $H_2(x; \sigma) = x^2 - \sigma$

(8)

We have

$$\int_{\mathbb{T}^2} :u_N^2(x): = \sum_{|n| \leq N} \frac{|g_n|^2 - 1}{\langle n \rangle^2}$$

- $\mathbb{E}\left(\int_{\mathbb{T}^2} :u_N^2: dx\right) = 0$
- $\mathbb{E}\left[\left(\dots\right)^2\right] = \mathbb{E}\left[\left(g_0^2 - 1\right)^2\right]$
 $+ 2\mathbb{E}\sum_{n \in I} \frac{\left(|g_n|^2 - 1\right)^2}{\langle n \rangle^4}$ "I = $\mathbb{Z}^2/2$ "
 $\lesssim 1 + \sum_{|n| \leq N} \frac{1}{\langle n \rangle^4} \lesssim 1 \text{ indep of } N$

$$\mathbb{E}\left[\left(|g_n|^2 - 1\right)\left(|g_m|^2 - 1\right)\right] = 0 \text{ unless } n = m.$$

$\Rightarrow \int_{\mathbb{T}^2} :u_N^2: dx$ is a well-defined r.v.

$$\int_{\mathbb{T}^2} :u^2: dx = \lim_{N \rightarrow \infty} \int :u_N^2: dx \text{ exists in } L^q(\mathbb{T})$$

for any $q < \infty$.