

# Phase dependence of escape from star clusters on elliptic orbits



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## Outline

In this short study the rate of escape from an idealised star cluster is investigated, by means of small N-body simulations. The cluster is subject to the tidal field of a galaxy modelled as a point mass. It is shown that the dependence of the lifetime on the initial orbital phase is negligible, even for an eccentricity  $e = 0.5$ . It is deduced that, for galactic orbits with a given semi-major axis, the lifetime is independent of  $e$  to first order in  $e$ . The results on the lifetime of a star cluster on an elliptical orbit agree with the formula given recently by Cai et al. (2015).

## Method and initial conditions

Simulations are run with 8192 stars of mass  $1M_{\odot}$ , no primordial binaries, and no stellar evolution. The initial model is Plummer's model with a virial radius of 1pc, but the model is truncated at ten scale radii. It is in orbit about a galaxy represented by a point mass of mass  $M_g = 10^{10}M_{\odot}$ , and the semi-major axis is 1211pc. For a cluster on a circular orbit, it follows that the initial model has  $r_h/r_J = 0.1$ , where  $r_h, r_J$  are the half-mass and Jacobi radii, respectively. For models on an elliptic galactic orbit the eccentricity is  $e = 0.5$ , and the initial phase (strictly, the initial value of the eccentric anomaly) is  $E_0 = 0, 90, 180$  or  $270^{\circ}$ .

Simulations were run with NBODY6 Nitadori & Aarseth (2012), in a non-rotating frame initially centred at the barycentre of the cluster. All stars were retained in the simulation. The bound mass  $M$  was defined to be the mass of stars within a sphere, centred at the density centre, with radius  $r$  such that

$$r = R_g \left( \frac{M}{3M_g} \right)^{1/3},$$

where  $R_g$  is the instantaneous distance from the galactic centre. The lifetime  $T_{0.1}$  is defined to be the time at which  $M$  is 0.1 times its initial value. In this poster quantities such as  $M$  and  $T_{0.1}$  are expressed in Hénon units (Hénon, 1971).

## Results of the N-body simulations – I

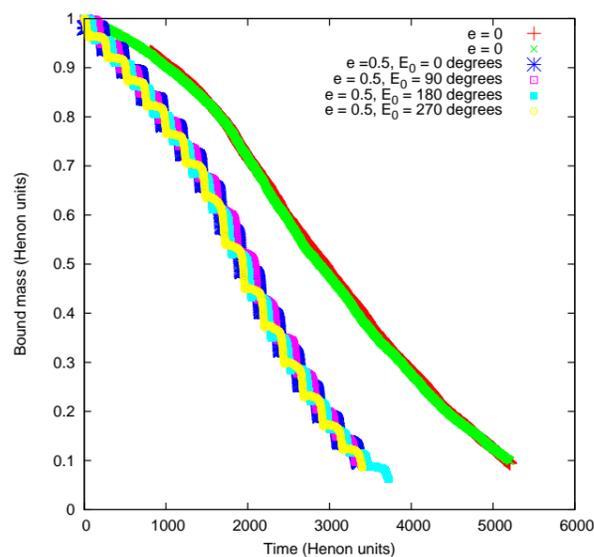


Fig.1 Bound mass against time for two models on circular orbits, and four on elliptic orbits in which the initial "phase" is  $0^{\circ}$  (perigalacticon),  $90, 180$  and  $270^{\circ}$ . Early data for one of the circular models is missing.

### Immediate Conclusions

1. The mass-loss rate is greater on an elliptical orbit than on a circular orbit with the same semi-major axis
2. For a given eccentricity and semi-major axis, the mass-loss rate is almost independent of the starting phase (i.e. starting at apogalacticon, perigalacticon, or points in between).

These conclusions are quantified in the next panel.

## Results of the N-body simulations – II

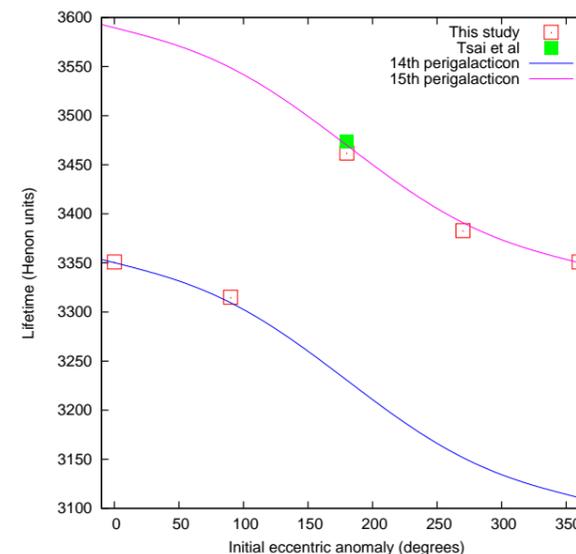


Fig.2 Lifetime  $T_{0.1}$  as a function of the initial "phase" (eccentric anomaly,  $E_0$ ). The curves give the time of the 14th and 15th perigalactic passage; note that the 15th becomes the 14th as  $E_0$  increases through  $0$  or  $360^{\circ}$ . Also shown is the result obtained from the formula for  $T_{0.1}(e)/T_{0.1}(0)$  given by Cai et al. (2015), where  $T_{0.1}(e)$  is the lifetime on an orbit of eccentricity  $e$  and given semi-major axis. It is plotted at  $E_0 = 180^{\circ}$ , which was their initial setup, and the value of  $T_{0.1}(0)$  was obtained from the average of the two models in Fig.1 of this poster.

### Immediate Conclusions

1. For the setup in this study the lifetime on an elliptical orbit of eccentricity 0.5 is 0.670 times the lifetime on a circular orbit with the same semi-major axis.
2. The lifetime is consistent with an escape rate which is independent of the initial "phase"  $E_0$ . The lifetime is reached close to the first perigalacticon where the remaining mass of the cluster falls below 0.1 times its initial value.

## Implications

### 1. Dependence on eccentricity

Let us suppose that the lifetime is independent of  $E_0$ , i.e.  $T_{0.1}(e, E_0) = T_{0.1}(e)$ . The equations of Kepler motion

$$x = a(1 - e \cos E), y = b \sin E$$

are invariant under the transformation

$$e \rightarrow -e, E \rightarrow \pi - E.$$

Thus motion on a galactic orbit of eccentricity  $-e$  is the same as motion on an orbit of eccentricity  $e$ , but with a different value of  $E_0$ . From the above assumption it follows that  $T_{0.1}(-e) = T_{0.1}(e)$ , i.e.  $T_{0.1}$  is an even function of  $e$ . This simplification was adopted in the fitting formula (for the dependence of  $T_{0.1}$  on  $e$ ) presented by Cai et al. (2015).

### 2. The role of perigalacticon

For small eccentricity  $e$ , it follows that  $T_{0.1}(e) = T_{0.1}(0) + O(e^2)$ . Thus the lifetime on a slightly elliptical orbit is almost the same as the lifetime on a circular orbit with the same semi-major axis. This is an orbit with radius almost equal to the average of the peri- and apo-galactic radii of the elliptical orbit. Thus mass loss from a cluster on a mildly elliptical orbit is not determined by conditions at perigalacticon. It is already known that the half-mass radius is not set at perigalacticon either (Küpper et al., 2010).

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## References

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