

6. Why I love the 3-body problem

1. "Star-studded" history

... the great names of dynamics

2. A clean problem

... the only mess is in one's head

3. A cross-roads in science

... where mathematics, physics and astronomy all meet.

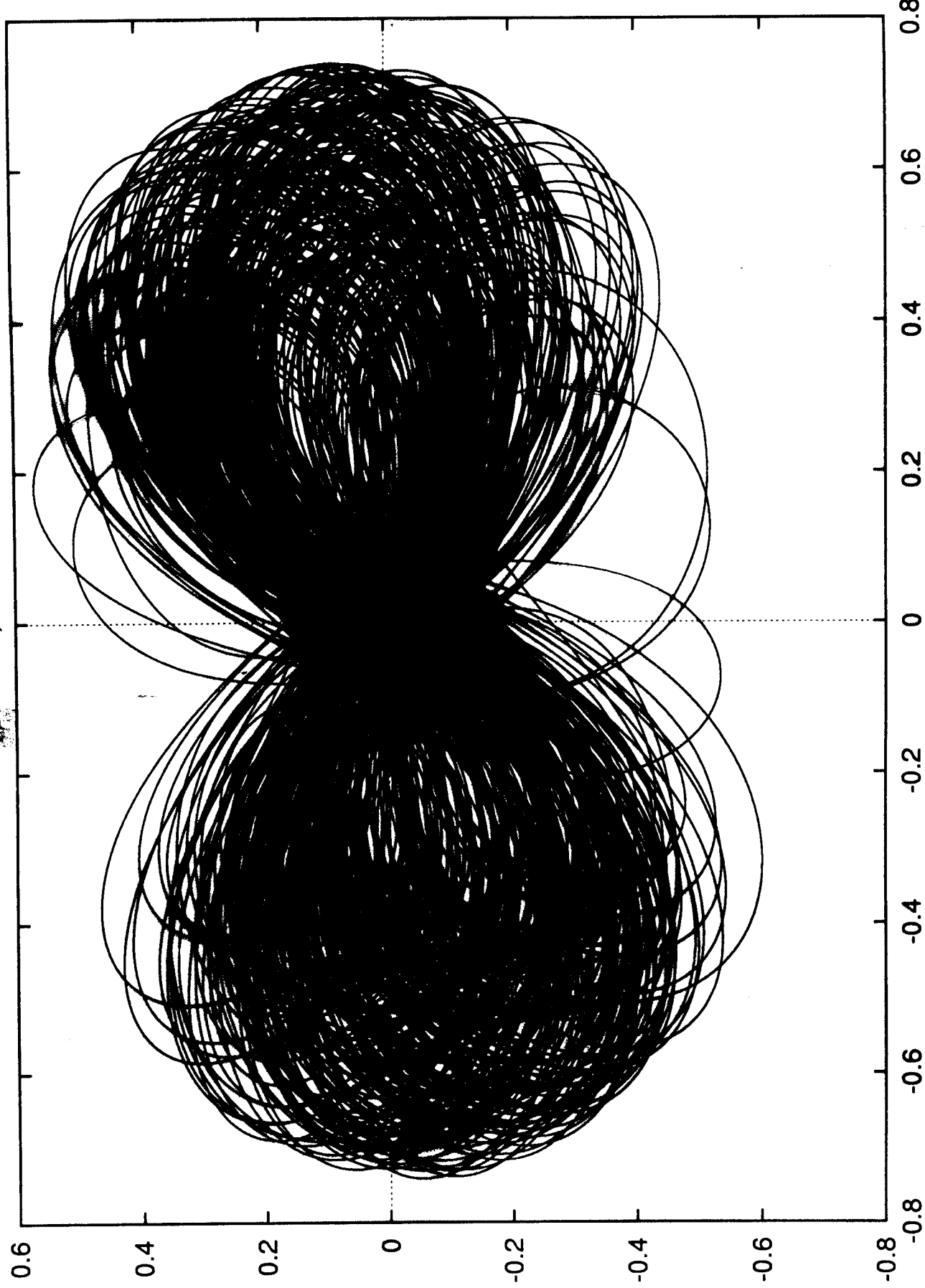
4. Self-rejuvenating

... new problems always emerging from new applications and new insights.

'... every generation formulates and solves its own "fundamental questions in the three-body problem".'

[Wintner]

Seed 18, Case 55



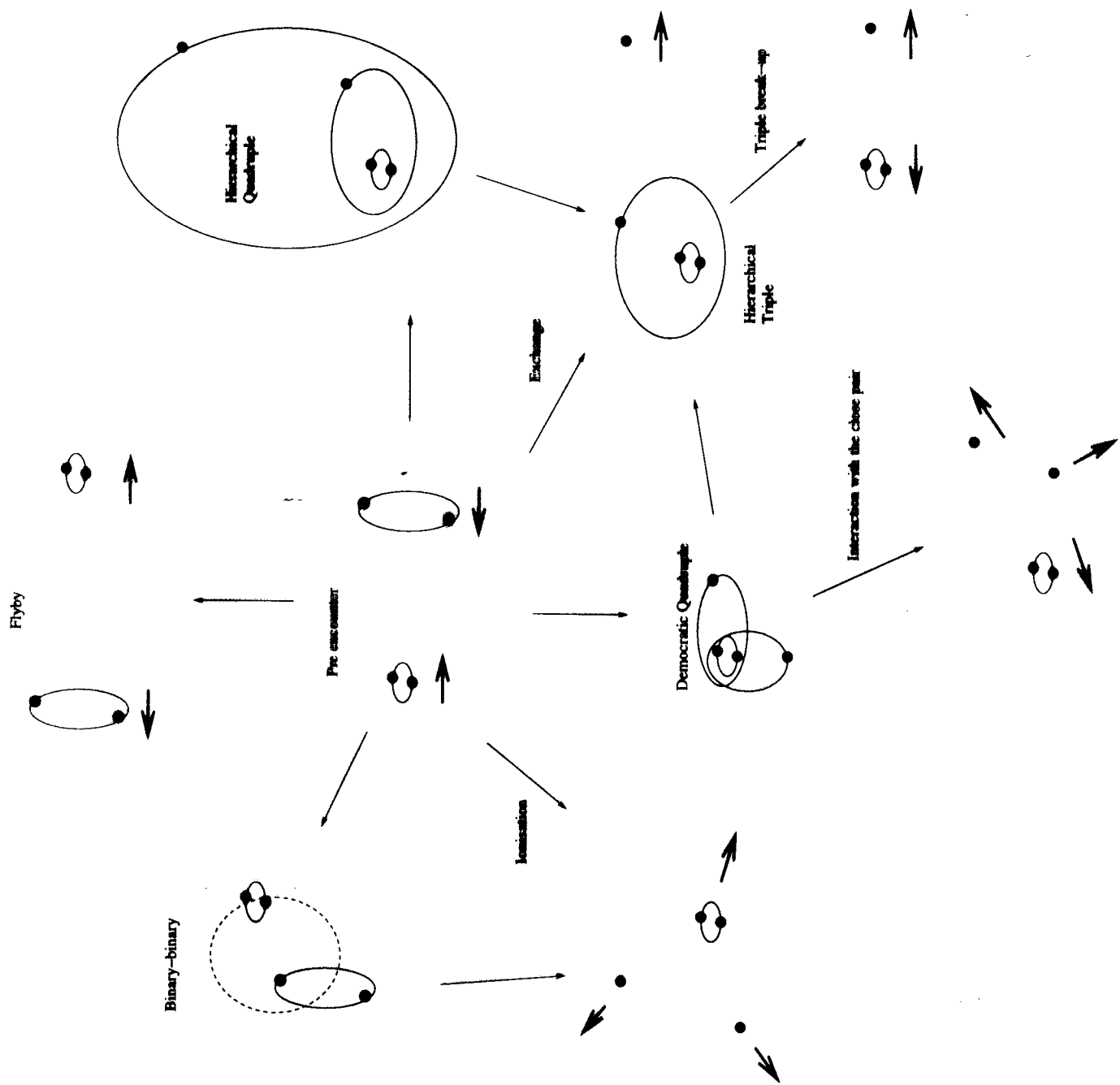
4. Proportion of 8-like systems

- (a) select x_i, y_i, u_i, v_i uniformly on $(0,1)$
- (b) shift to barycentre
- (c) rotate to make $\sum m_i x_i v_i = 0$
- (d) scale to virial equilibrium $2T - U = 0$.
- (e) scale to standard units: $M = G = 4T = 2U = 1$.

Result

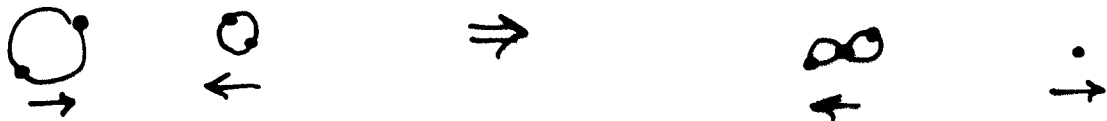
- (i) $\approx 5\%$ 8-like to $t=10$
- (ii) $\lesssim 1\%$ 8-like to $t=100$
- (iii) $\approx 0.1\%$ 8-like to $t=1000$

Conclusion: quite an improbable outcome of binary-binary scattering



Binary- Binary Scattering

1. Outcome analysed in terms of asymptotic solutions
2. New possible outcome

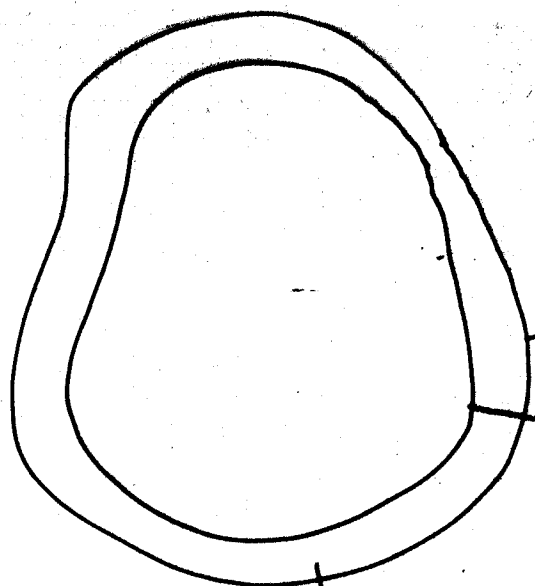


3. Recognising an 8-like orbit automatically

- (i) find change in orientation, i.e. sign change in $(\underline{r}_2 - \underline{r}_1) \times (\underline{r}_3 - \underline{r}_1)$, to detect collinearity
- (ii) determine central body at each collinear config
- (iii) look for sequence ... 123123123... or reverse (Such a system is described as "8-like").

The impossibility of permanent capture

Let R be some region in phase space with finite volume, e.g. $|\underline{r}_i - \bar{r}| < L, i=1,2,J$



I: all $x(0)$ s.t. $x(t) \in R \forall t \geq 0$

II: all $x(1)$ s.t. $x(t) \in R \forall t \geq 0$

= all $x(0)$ s.t. $x(t) \in R \forall t \geq -1$

$\subset$ all $x(0)$ s.t. $x(t) \in R \forall t \geq 0$

III: 0 volume -

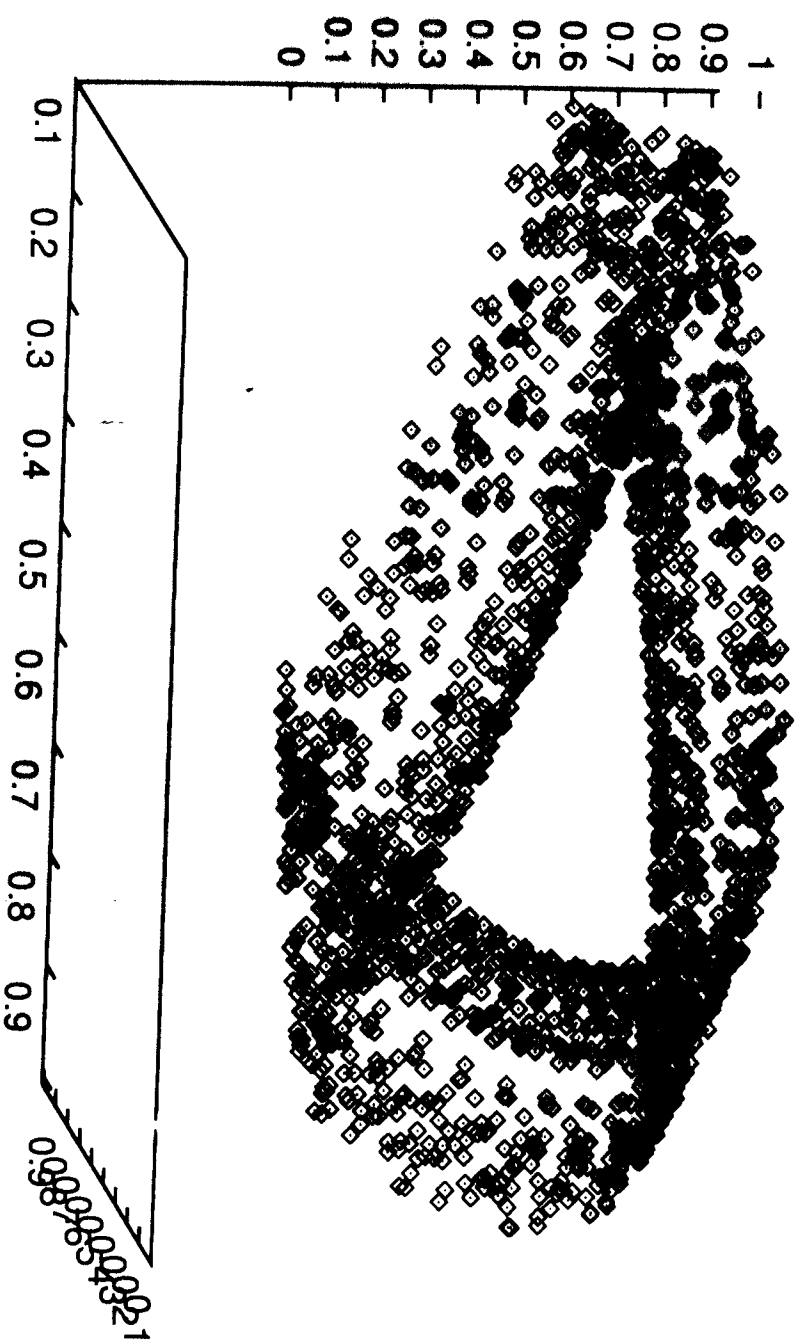
all $x(0)$ s.t. $x(t) \in R \forall t \geq 0$

& $x(t) \notin R$ for some $t \in [-1, 0]$

= all $x(0)$ s.t. $x(t)$ becomes trapped in $[-1, 0]$

Displaced central particle

'triple.out' $\diamond$



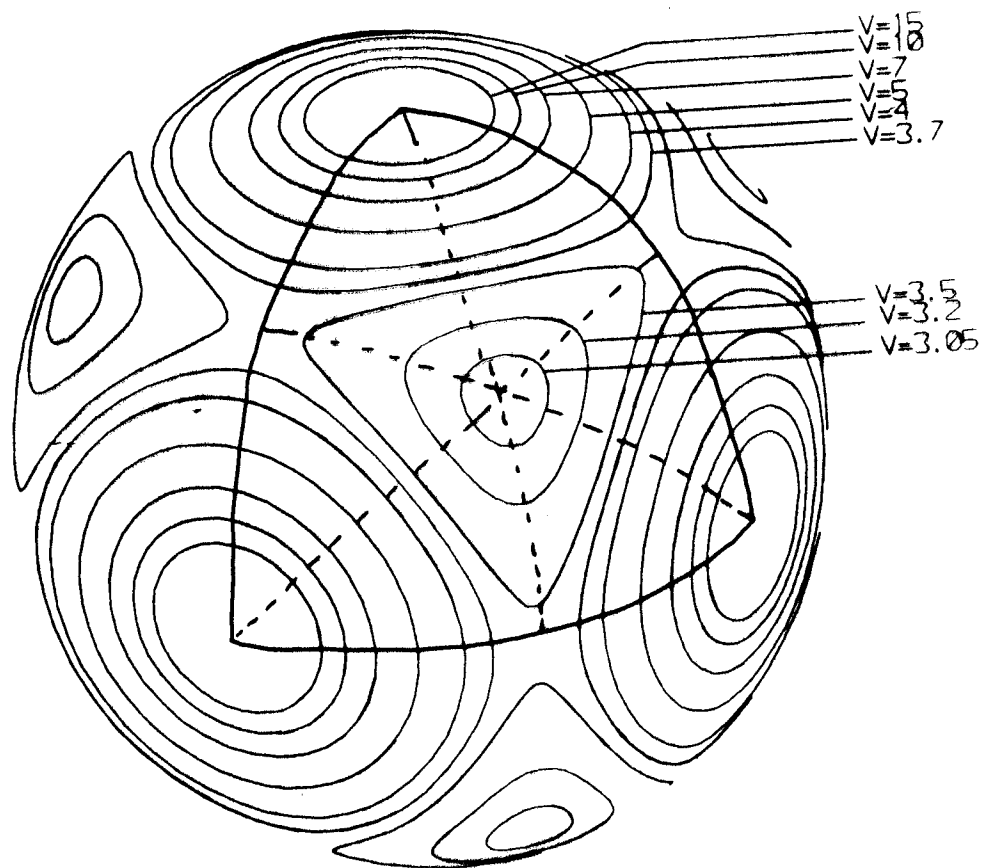


Fig. 4. The surface $\bar{I} = \text{const}$ of Equation (58) in configuration space $(\bar{\alpha}_1, \bar{\alpha}_2, \bar{\alpha}_3)$. The lines are contours $\bar{O} = \text{const}$. Equal mass case $m_1 = m_2 = m_3$.

2. Exclude collisions

(a) Suppose collision between m_2 & m_3 . Then

$$\begin{aligned} \int (T_{\text{def}} + U) dt &> \int (T_{\text{def}}^* + U^*) dt \\ &\quad \uparrow \qquad \qquad \uparrow \\ &\quad T|_{m_1=0} \quad U|_{m_1=0} \\ &\geq A_2 \equiv \inf_{\substack{\text{2-body} \\ \text{collision orbits}}} \int (T_{\text{def}}^* + U^*) dt \end{aligned}$$

(b) Construct non-collision path C from $L_3 \cap I_3$ to I_1 along contour of constant U , described at constant speed; length $= \ell_0$.

$$\text{Then } \inf \int (T_{\text{def}} + U) dt < \int_C (T_{\text{def}} + U) dt < A_2$$

provided that $\ell_0 < \frac{\pi}{5}$.

Numerical quadrature shows

$$\ell_0 < \frac{\pi}{5.08255\dots}$$

3. Show that minimiser joins smoothly

(to next segment from I_1 to $L_2 \cap I_2$)

Existence of Segment from $L_3 \cap I_3$ to I_1

1. Extremise action $\int \mathcal{L} dt$

$$\text{where } \mathcal{L} = T + U \quad \leftarrow \sum_{i < j} \frac{m}{|\mathbf{r}_i - \mathbf{r}_j|}$$

$\uparrow$
kinetic energy

$$(a) \quad T = T_{\text{rot}} + T_{\text{def}} \quad \leftarrow \text{k.e. of deformation}$$

$\uparrow$

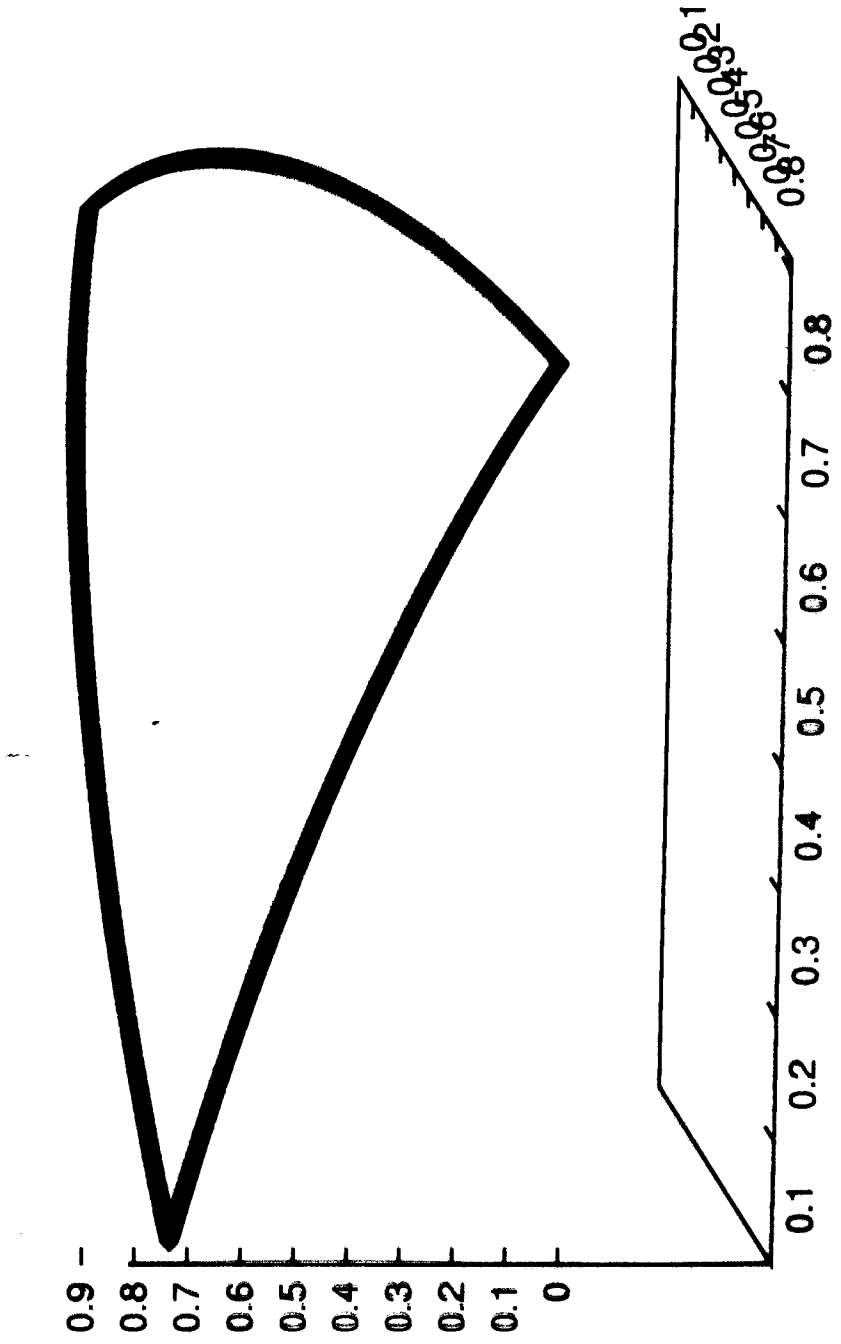
$$\text{k.e. of rotation} = \frac{C^2}{I} \quad \leftarrow \begin{array}{l} \text{angular momentum} \\ \text{moment of inertia} \end{array}$$

attempt to minimise $\int \mathcal{L} dt$ by assuming $C=0$,
i.e. $\delta \int (T_{\text{def}} + U) dt = 0$.

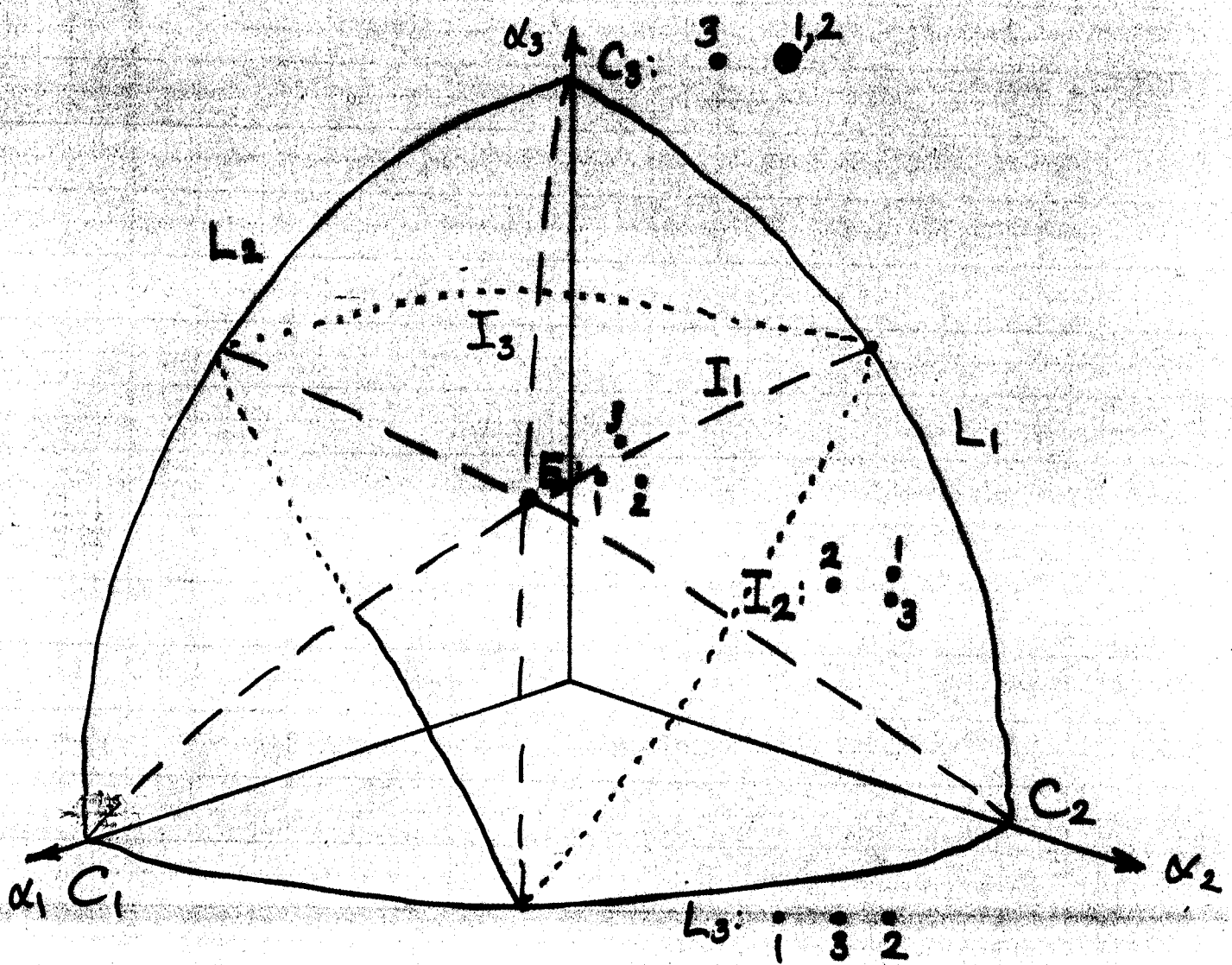
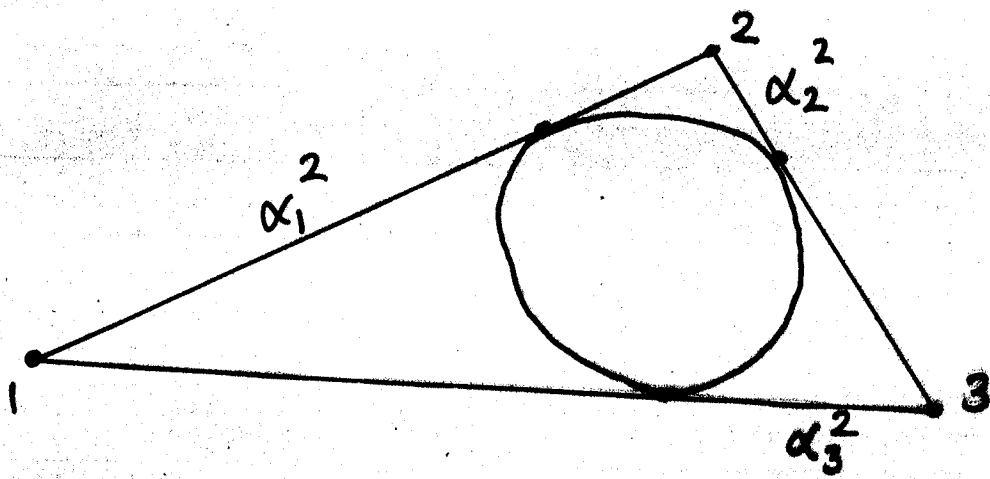
Existence of minimiser "standard".

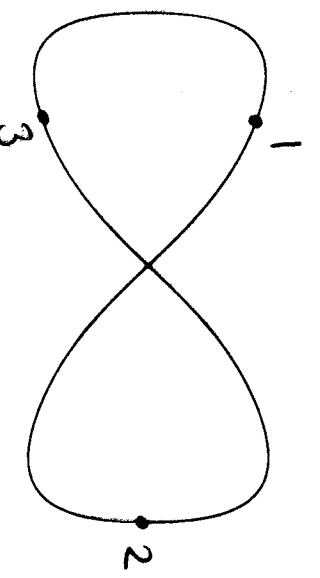
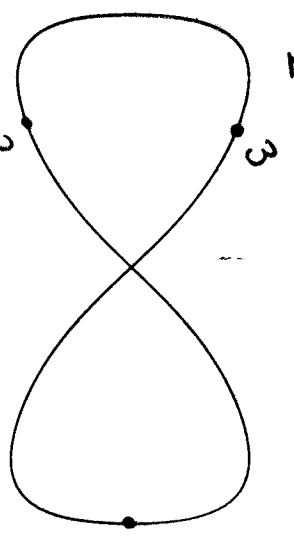
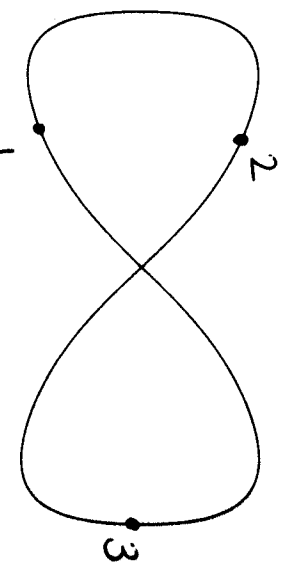
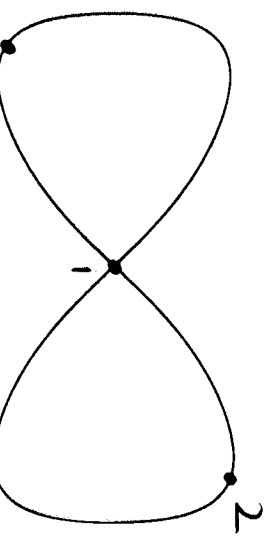
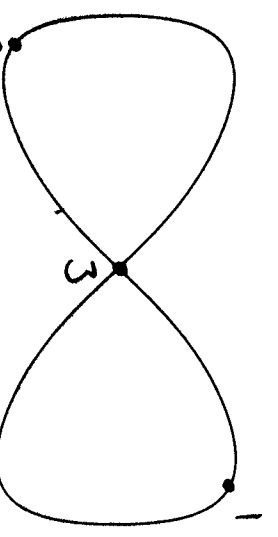
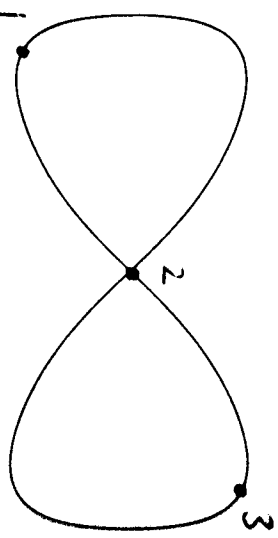
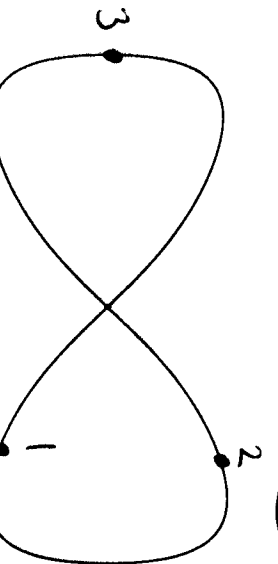
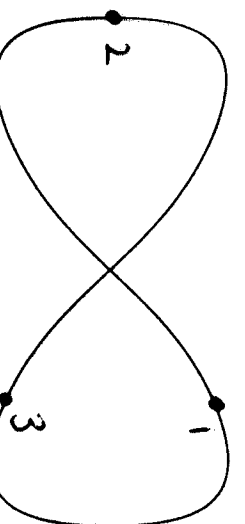
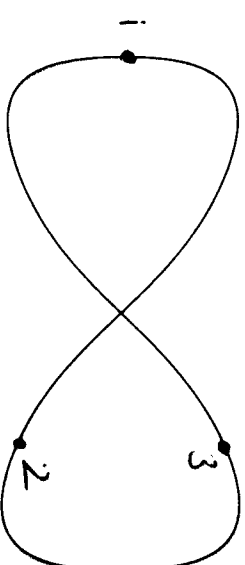
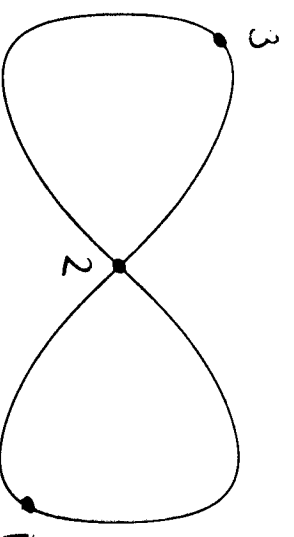
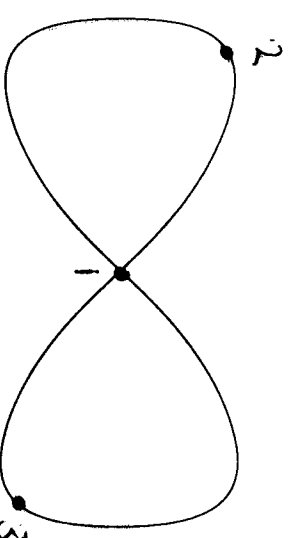
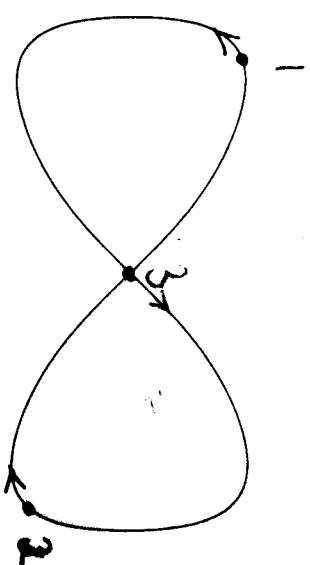
Figure 8

'triple.out' ◇



Lemaitre's Parametrisation





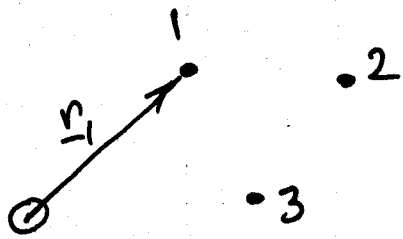
↑

↑

↑

↑

Classical Gravitational 3-Body Problem



$$\ddot{\underline{r}}_i = -G \sum_{\substack{j=1 \\ j \neq i}}^3 m_j \frac{\underline{r}_i - \underline{r}_j}{|\underline{r}_i - \underline{r}_j|^3}, \quad i=1,2,3$$

Take $Gm_i = 1$ (scaled, equal masses)

$\underline{r}_i \in \mathbb{R}^2$ (planar)

At $t=0$

$$\underline{r}_1 = (a, b), \quad a = -.97000436 \\ b = -.24308753$$

$$\underline{r}_2 = -\underline{r}_1$$

$$\underline{r}_3 = \underline{0}$$

$$\dot{\underline{r}}_3 = (c, d), \quad c = 0.93240737 \\ d = 0.86473146$$

Source <http://orca.ucsc.edu/~rmont/index.html>